

# ENERGY SUBCRITICAL NLS ON $\mathbf{R}^N \times \mathbf{R}/2\pi\mathbf{Z}$

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## 1. RESUME

The main object of the talk is to present some results concerning the non-linear Schrödinger equation (NLS) on the product space  $\mathbf{R}^n \times \mathbf{T}$ , with  $n \geq 1$  and  $\mathbf{T} = \mathbf{R}/2\pi\mathbf{Z}$  is the one-dimensional standard torus. More precisely we consider the following family of Cauchy problems:

$$(1.1) \quad \begin{cases} i\partial_t u + \Delta_{x,y} u + \lambda |u|^\alpha u = 0, & (t, x, y) \in \mathbf{R} \times \mathbf{R}^n \times \mathbf{T} \\ u(0, x, y) = f(x, y) \in H_{x,y}^1, \end{cases}$$

where

$$\Delta_{x,y} = \sum_{l=1}^n \partial_{x_l}^2 + \partial_y^2,$$

$\lambda$  is a real number, the nonlinearity parameter  $\alpha$  satisfies the assumption

$$\alpha < \alpha^*(n), \quad \alpha^*(n) = \begin{cases} \frac{4}{n-1} & \text{if } n \geq 2, \\ +\infty & \text{if } n = 1. \end{cases}$$

and

$$H_{x,y}^1 = (1 - \Delta_{x,y})^{-\frac{1}{2}} L_{x,y}^2, \text{ where } L_{x,y}^2 = L^2(\mathbf{R}^n \times \mathbf{T}).$$

We underline that the power nonlinearity  $\alpha^*(n)$  corresponds to the  $H^1$ -critical NLS in  $\mathbf{R}^{n+1}$  as well as to the  $H^{1/2}$ -critical nonlinearity in  $\mathbf{R}^n$ .

In this regime we are able to obtain the following:

- (1) for any initial datum  $f \in H_{x,y}^1$ , the problem (1.1) has a unique local solution

$$u(t, x, y) \in C((-T, T); H_{x,y}^1),$$

where  $T = T(\|f\|_{H_{x,y}^1}) > 0$ ;

- (2) there exists  $\varepsilon_0 > 0$  such that the flow

$$B_{\varepsilon_0}(0) \cap H_{x,y}^1 \ni f \rightarrow u(t, x, y) \in C((-T, T); H_{x,y}^1),$$

is Lipschitz continuous.

Moreover, by using a conservation laws argument, one can see that the local solutions above are global in time in the following cases:

- $\lambda < 0$  (i.e. (1.1) is defocusing);
- $\lambda \in \mathbf{R}$  and  $0 < \alpha < 4/(n+1)$  (i.e. (1.1) is  $L^2$ -subcritical);

- $\lambda > 0$ ,  $4/(n+1) \leq \lambda < 4/(n-1)$  and initial data  $f$  small enough in  $H_{x,y}^1$  (i.e. (1.1) is energy-subcritical,  $L^2$ -supercritical, focusing and the initial data are small).

We recall that the key ingredient used in the literature in order to get (at least) local well-posedness for (1.1) is a suitable version of Strichartz estimates of the type

$$\|e^{it\Delta_{x,y}} f\|_{L_t^q L_x^r L_y^2} \leq C \|f\|_{L_{x,y}^2},$$

where  $(q, r)$  are such that

$$(1.2) \quad \frac{2}{q} + \frac{n}{r} = \frac{n}{2}, \quad q \geq 2, \quad (n, q) \neq (2, 2).$$

On the other hand it is well-known (by classical Euclidean theory) that the best nonlinearity that can be reached with this technique is the  $L^2$ -critical nonlinearity in  $\mathbf{R}^n$ , i.e.  $0 < \alpha < 4/n$ . Our main contribution is the analysis of (1.1) in the case  $4/n \leq \alpha < \alpha^*(n)$  (i.e. nonlinearities which are  $L^2$ -supercritical and  $H^{1/2}$ -subcritical in  $\mathbf{R}^n$ ). The main difficulty in the transposition of the analysis above in this larger regime of nonlinearities is that, in analogy with the theory of  $L^2$ -supercritical NLS in  $\mathbf{R}^n$ , it seems to be necessary to work with Strichartz estimates involving derivatives w.r.t.  $x$  variable.

To overcome this obstacle we exploit a class of inhomogeneous Strichartz estimates with respect the  $x$  variable true in a range of Lebesgue exponents larger than the one given in the usual homogeneous estimates context. This allow us to perform a fixed point argument by considering in our partially periodic setting the same numerology involved in the analysis of NLS posed in the euclidean space  $\mathbf{R}^n$  via admissible Strichartz norms  $L_t^q L_x^r$ , with  $(q, r)$  as in (1.2).