

Provided for non-commercial research and educational use.
Not for reproduction, distribution or commercial use.

Mathematica Balkanica

Mathematical Society of South-Eastern Europe
A quarterly published by
the Bulgarian Academy of Sciences – National Committee for Mathematics

The attached copy is furnished for non-commercial research and education use only. Authors are permitted to post this version of the article to their personal websites or institutional repositories and to share with other researchers in the form of electronic reprints.

Other uses, including reproduction and distribution, or selling or licensing copies, or posting to third party websites are prohibited.

For further information on Mathematica Balkanica visit the website of the journal
<http://www.mathbalkanica.info>

or contact:

Mathematica Balkanica - Editorial Office;
Acad. G. Bonchev str., Bl. 25A, 1113 Sofia, Bulgaria
Phone: +359-2-979-6311, Fax: +359-2-870-7273,
E-mail: balmat@bas.bg

Limited Domain Radon Transform ¹

Árpád Kurusa

Presented by P. Kenderov

The problem in this article is to recover a function on \mathbb{R}^n from its integrals known only on hyperplanes intersecting

1. Introduction

There is a number of papers concerning the reconstruction of a function from only a partial knowledge of the function's Radon transform. The two most known examples are the exterior Radon transform [6] and the limited angle Radon transform [4].

The classical Radon transform is defined for an integrable function f on \mathbb{R}^n by

$$Rf(\omega, p) = \int_{H(\omega, p)} f(x) dx_h,$$

where $\omega \in S^{n-1}$ is a unit vector, $p \in \mathbb{R}_+$ and $Rf(\omega, p)$ is just the integral of f over the hyperplane $H(\omega, p) = \{x \in \mathbb{R}^n : \langle x, \omega \rangle = p\}$ by the surface measure dx_h on it.

In the limited angle case $Rf(\omega, p)$ is restricted in ω to a subset of S^{n-1} . The exterior Radon transform is the restriction of Rf to the set $p > 1$.

We define the *limited domain Radon transform* $R_L f$ of a function as the restriction of Rf onto the set $p \leq 1$. In the next section we prove its continuity on a weighted class $L^2_{\alpha, \beta}(\mathbb{R}^n)$ of square integrable functions that are zero in a neighbourhood of the origin. In Section 3 we give the null space and range of R_L acting on $L^2_{\alpha, \beta}(\mathbb{R}^n)$ for odd dimensional spaces. In Section 4 we do the same for even dimensional spaces, where the injectivity of R_L on $L^2_{\alpha, \beta}(\mathbb{R}^n)$ turns out.

¹Supported by the Hungarian NFS, Grants No. T4427, W015452, F016226 and T020066

The author thanks the Mathematisches Institut of the Universität der Erlangen-Nürnberg and the Department of the University of Maryland for their support and assistance while working on this paper.

2. Preliminaries

The limited domain Radon transform of a function f integrable on the hyperplanes is defined by means of the formula $R_L f = \chi_{[0,1)} Rf$, where $\chi_{[0,1)}$ is the characteristic function of the interval $[0, 1)$.

From the well known inversion formulas [3] we know that a square integrable function is determined by the limited domain Radon transform in the unit open ball B^n , if the dimension n is odd. In even dimensions the situation is more complicated, but still there is an approximate possibility, as C. Berenstein and D. Walnut proved in [1] that to recover f to a given accuracy in a ball of radius $R > 0$ it is sufficient to know $Rf(\omega, p)$ only for $p < R + \alpha$ for some $\alpha > 0$ that depends on the accuracy desired (the greater the accuracy desired, the greater α must be).

Motivated by these observations, in the next two sections we carry our investigations onto the following spaces. $L^2_{\alpha,\beta}(\mathbb{E}^n)$ is the Hilbert space of functions supported in $\mathbb{E}^n = \mathbb{R}^n \setminus \text{Cl } B^n$ equipped with the inner product

$$\langle f, g \rangle_{\alpha,\beta} = \int_{\mathbb{E}^n} f(x)g(x)|x|^\alpha(1-|x|^{-2})^\beta dx.$$

Estimate of the weight shows immediately, that $L^2_{\alpha,\beta}(\mathbb{E}^n) \supseteq L^2_{\alpha',\beta'}(\mathbb{E}^n)$ for $\alpha \leq \alpha'$ and $\beta \leq \beta'$. Denote $\mathbb{B}^n = \{H(\omega, p) : 0 \leq p \leq 1, \omega \in S^{n-1}\}$. Then the Hilbert space of functions on \mathbb{B}^n with the inner product

$$\langle F, G \rangle_{\gamma,\delta} = \int_{S^{n-1}} \int_0^1 F(\omega, p)G(\omega, p)p^\gamma(1-p^2)^\delta dp d\omega$$

is $L^2_{\gamma,\delta}(\mathbb{B}^n)$. Here $L^2_{\gamma,\delta}(\mathbb{B}^n) \supseteq L^2_{\gamma',\delta'}([0, 1])$ for $\gamma' \leq \gamma$ and $\delta' \leq \delta$. Finally we need the Hilbert space $L^2_{\gamma,\delta}([0, 1])$ of functions on $[0, 1]$ with the inner product

$$\langle f, g \rangle_{\gamma,\delta} = \int_0^1 f(p)g(p)p^\gamma(1-p^2)^\delta dp.$$

For $-1 < \gamma \leq 0$, $-1 < \delta \leq 0$, the polynomials are dense in this space, and $L^2_{\gamma,\delta}([0, 1]) \supseteq L^2_{\gamma',\delta'}([0, 1])$ for $\gamma' \leq \gamma$ and $\delta' \leq \delta$.

Our main tool is the spherical harmonic expansion in these spaces. Briefly, the spherical harmonics, $Y_{\ell,m}$ constitute a complete polynomial orthonormal system in the Hilbert space $L^2(S^{n-1})$. If $f \in C^\infty(S^{n-1} \times \mathbb{R}_+)$ and $f_{\ell,m}(p)$ is the

corresponding coefficient of $Y_{\ell,m}(\omega)$ in the expansion of $f(\omega, p)$, i.e. $f_{\ell,m}(p) = \int_{S^{n-1}} f(\omega, p) \overline{Y_{\ell,m}(\omega)} d\omega$, then the series $\sum_{\ell,m}^{\infty} f_{\ell,m}(p) Y_{\ell,m}(\omega)$ converges uniformly absolutely on compact subsets of $S^{n-1} \times \mathbb{R}$ to $f(\omega, p)$. For further references we refer to [7]. Below we shall use the expansions

$$g(\varphi, p) = \sum_{m=-\infty}^{\infty} g_m(p) \exp(im\varphi) \quad \text{and} \quad g(\omega, p) = \sum_{\ell,m}^{\infty} g_{\ell,m}(p) Y_{\ell,m}(\omega)$$

in dimension two and in higher dimensions, respectively. In dimension two, φ will mean the angle of the respective unit vector to a fixed direction.

The spherical expansions of the Radon transforms are well known [6]. Applying these to the functions in $L^2_{\alpha,\beta}(\mathbb{E}^n)$ we obtain

$$(2.1) \quad (R_L f)_m(p) = 2 \int_1^{\infty} f_m(q) \frac{\cos(m \arccos(p/q))}{\sqrt{1-p^2/q^2}} dq$$

for dimension two and

$$(2.2) \quad (Rf)_{\ell,m}(p) = \frac{|S^{n-2}|}{C_m^{\lambda}(1)} \int_1^{\infty} f_{\ell,m}(q) q^{n-2} C_m^{\lambda}\left(\frac{p}{q}\right) \left(1 - \frac{p^2}{q^2}\right)^{\frac{n-3}{2}} dq$$

for higher dimensions, where C_m^{λ} is the Gegenbauer polynomial of degree m , $\lambda = (n-2)/2$ and $p \leq 1$.

An important consequence of these expansions is the continuity of R_L .

Theorem 2.1. R_L maps $L^2_{\alpha,\beta}(\mathbb{E}^n)$ continuously into $L^2_{\gamma,\delta}(\mathbb{B}^n)$, where $\alpha > n-2$, $\beta < 1$, $\gamma > -1$ and $\delta > \begin{cases} -1 & \text{if } n \geq 3 \\ -1/2 & \text{if } n = 2 \end{cases}$.

Proof. First, observe that

$$(2.3) \quad \|R_L f\|_{\gamma,\delta}^2 = \sum_{\ell,m} \|Y_{\ell,m}\|_2^2 \int_0^1 (R_L f)_{\ell,m}^2(p) p^{\gamma} (1-p^2)^{\delta} dp.$$

Using (2.1), (2.2) and that $|C_m^{\lambda}(x)| \leq |C_m^{\lambda}(1)|$ for $|x| \leq 1$ we can estimate (2.3) by

$$c_1 \int_0^1 \left| \int_1^{\infty} f_{\ell,m}(q) q^{n-2} \left(1 - \frac{p^2}{q^2}\right)^{\frac{n-3}{2}} dq \right|^2 p^{\gamma} (1-p^2)^{\delta} dp,$$

where c_1 is a suitable constant independent from m . For $n \geq 3$, this is less than

$$(2.4) \quad c_1 \left| \int_1^{\infty} f_{\ell,m}(q) q^{n-2} dq \right|^2 \int_0^1 p^{\gamma} (1-p^2)^{\delta} dp.$$

For $n = 2$ we estimate by $|1 - p^2/q^2| \geq 1 - p^2$ and obtain

$$(2.5) \quad c_1 \left| \int_1^\infty f_{\ell,m}(q) dq \right|^2 \int_0^1 p^\gamma (1 - p^2)^{\delta-1/2} dp.$$

The second integrals in (2.4) and (2.5), respectively, explain the restrictions on γ and δ given in the theorem.

The restrictions on α and β come from the following estimate of the first integrals in (2.4) and (2.5).

$$(2.6) \quad \begin{aligned} \left| \int_1^\infty f_{\ell,m}(q) q^{n-2} dq \right|^2 &= \left| \int_1^\infty f_{\ell,m}(q) q^{-1-\alpha} (1 - q^{-2})^{-\beta} q^{\alpha+n-1} (1 - q^{-2})^\beta dq \right|^2 \\ &\leq \int_1^\infty f_{\ell,m}^2(q) q^{\alpha+n-1} (1 - q^{-2})^\beta dq \\ &\quad \times \int_1^\infty q^{-2-2\alpha} (1 - q^{-2})^{-2\beta} q^{\alpha+n-1} (1 - q^{-2})^\beta dq. \end{aligned}$$

The last two integrals need to be finite to ensure the finiteness of (2.3). The first one is finite because $f \in L_{\alpha,\beta}(\mathbb{E}^n)$, the other is finite if $\alpha > n - 2$ and $\beta < 1$. ■

Throughout this paper we shall assume that α , β , γ and δ satisfy the conditions given in Theorem 2.1.

The weight $(1 - p^2/q^2)^{(n-3)/2}$ is substantially different in odd and in even dimensions; it is polynomial in odd dimensions.

3. Odd dimensions

Let $n = 2d + 3$ and $d \geq 0$. Then $\lambda = d + 1/2$ and

$$(3.1) \quad (R_L f)_{l,m}(p) = \frac{|S^{n-2}|}{C_m^\lambda(1)} \int_1^\infty f_{\ell,m}(q) q^{n-2} D_m^\lambda \left(\frac{p}{q} \right) dq,$$

where $D_m^\lambda(x) = C_m^\lambda(x)(1 - x^2)^d$.

Lemma 3.1. D_m^λ is a polynomial of the form

$$D_m^\lambda(x) = \sum_{i=0}^{d+[m/2]} d_{m,i}^\lambda x^{m+2d-2i},$$

where the coefficients $d_{m,i}^\lambda \neq 0$.

The proof is a simple consequence of the formulas (8.932.1), (9.136.1) and (9.131.1) in [2].

Substituting the polynomial form of D_m^λ into (3.1) we see that $(R_L f)_{l,m}$ is also a polynomial of the form

$$(3.2) \quad (R_L f)_{l,m}(p) = \frac{|S^{n-2}|}{C_m^\lambda(1)} \sum_{i=0}^{d+[m/2]} p^{m+2d-2i} d_{m,i}^\lambda \int_1^\infty f_{\ell,m}(q) q^{1-m+2i} dq.$$

Let $2\phi_{\ell,m}(x^2) = f_{\ell,m}(1/x)/x^n$ and change the variable $q = 1/\sqrt{x}$. Then the integral in (3.2) becomes

$$(3.3) \quad c_{\ell,m,i} = \int_1^\infty f_{\ell,m}(q)q^{1-m+2i}dq = \int_0^1 \phi_{\ell,m}(x)x^{d-i+m/2} \frac{dx}{\sqrt{x}}.$$

Theorem 3.2. *The null space of R_L in $L^2_{\alpha,\beta}(\mathbb{E}^n)$ is the closure of the span of the functions*

$$g_{k,\ell,m}(\omega, q) = P_k^{(0,\varepsilon_m)}(2q^{-2} - 1)q^{-n}Y_{\ell,m}(\omega),$$

where $P_k^{(\cdot)}$ is the Jacobi polynomial of degree k , $n-1 \leq \alpha < n$, $-1 < \beta \leq 0$,

$$k \geq \left\lceil \frac{m+n-1}{2} \right\rceil \quad \text{and} \quad \varepsilon_m = \begin{cases} -1/2 & \text{if } m \text{ even} \\ 0 & \text{if } m \text{ odd} \end{cases}.$$

Proof. The fact that $g_{k,\ell,m} \in L^2_{\alpha,\beta}(\mathbb{E}^n)$ follows from $\alpha < n$ and $-1 < \beta$. Since also $n-1 \leq \alpha$ and $\beta \leq 0$, it follows $\phi_{\ell,m} \in L^2_{-1/2,0}([0,1])$.

According to (3.2) and (3.3) the function $R_L f$ vanishes, if and only if $c_{\ell,m,i} = 0$ for all ℓ, m , and $0 \leq i \leq d + [m/2]$.

For m even this gives that

$$(3.4) \quad 0 = \int_0^1 \phi_{\ell,m}(x)x^j \frac{dx}{\sqrt{x}} \quad \text{for all } 0 \leq j \leq d + m/2.$$

The shifted Jacobi polynomials $P_k^{(0,-1/2)}(2x-1)$ constitute a complete orthogonal system on $[0,1]$ with respect to the weight $1/\sqrt{x}$ [2(8.904)], therefore $\phi_{\ell,m}$ must be in the closure of the span of $\{P_k^{(0,-1/2)}(2x-1)\}_{k=d+1+[m/2]}^\infty$.

For m odd we have

$$(3.5) \quad 0 = \int_0^1 \phi_{\ell,m}(x)x^j dx \quad \text{for all } 0 \leq j \leq d + (m-1)/2.$$

The Jacobi polynomials $P_k^{(0,0)}(2x-1)$ constitute a complete orthogonal system on $[0,1]$ [2(8.904)], so $\phi_{\ell,m}$ is in the closure of the span of $\{P_k^{(0,0)}(2x-1)\}_{k=d+1+[m/2]}^\infty$.

The results of (3.4) and (3.5) give the theorem. ■

In the following we determine the range of the limited domain Radon transform.

Theorem 3.3. R_L maps the $L^2_{\alpha,\beta}(\mathbb{I}^n)$ closure of the span of the functions

$$h_{k,\ell,m}(\omega, q) = q^{-n-2k} Y_{\ell,m}(\omega), \quad 0 \leq k \leq d + [m/2]$$

where $n-1 \leq \alpha < n$ and $-1 < \beta < 0$, onto the $L^2_{\gamma,\delta}(\mathbb{I}^n)$ closure of the span of the functions

$$F_{\ell,m}(\omega, p) = Y_{\ell,m}(\omega) \sum_{i=0}^{d+[m/2]} p^{m+2d-2i} b_{\ell,m,i}$$

continuously and bijectively.

Proof. The easy verification of $h_{k,\ell,m} \in L^2_{\alpha,\beta}(\mathbb{I}^n)$ and $F_{\ell,m} \in L^2_{\gamma,\delta}(\mathbb{I}^n)$ is left to the reader. ■

Since $R_L: L^2_{\alpha,\beta}(\mathbb{I}^n) \rightarrow L^2_{\gamma,\delta}(\mathbb{I}^n)$ is continuous by Theorem 2.1, it maps a closed set into a closed set. Further, it is injective on the given functions by Theorem 3.2, therefore we only have to give coefficients $e'_{k,\ell,m} \in \mathbb{R}$ so that

$$F_{\ell,m} = R_L(f_{\ell,m} Y_{\ell,m}), \quad \text{where} \quad f_{\ell,m} Y_{\ell,m} = \sum_{k=0}^{d+[m/2]} e'_{k,\ell,m} h_{k,\ell,m}.$$

Eliminating $Y_{\ell,m}$ and reordering the summation we can search for $f_{\ell,m}$ in the form

$$(3.6) \quad f_{\ell,m}(q) = \sum_{k=0}^{d+[m/2]} e_{k,\ell,m} q^{-n} P_k^{(0,\varepsilon_m)}(2q^{-2} - 1),$$

where the coefficients $e_{k,\ell,m}$ are to be determined.

In $F_{\ell,m} = R_L(f_{\ell,m} Y_{\ell,m})$ the equality of two polynomials appears, which is equivalent to the coincidence of their corresponding coefficients. By (3.2) and (3.3) this gives

$$(3.7) \quad b_{\ell,m,i} = \frac{|S^{n-2}|}{C_m^\lambda(1)} d_{m,i}^\lambda \int_1^\infty f_{\ell,m}(q) q^{1-m+2i} dq$$

for $0 \leq i \leq d + [m/2]$. Substituting $f_{\ell,m}$ according to (3.6), we get a system of linear equations for $e_{k,\ell,m}$ with coefficients

$$\begin{aligned} a_{i,k,m} &= \int_1^\infty q^{1+2i-m-n} P_k^{(0,\varepsilon_m)}(2q^{-2} - 1) dq \\ &= \frac{1}{2} \int_0^1 x^{d+[m/2]-i} P_k^{(0,\varepsilon_m)}(2x - 1) x^{\varepsilon_m} dx, \end{aligned}$$

where $0 \leq k \leq d + [m/2]$. The Jacobi polynomials $P_k^{(0,\varepsilon_m)}(2x - 1)$ constitute orthogonal polynomial system with respect to the weight x^{ε_m} on $[0, 1]$, therefore

$a_{i,k,m} = 0$ for $k + i > d + [m/2]$ and $a_{i,d+[m/2]-i,m} \neq 0$ for $0 \leq i \leq d + [m/2]$. By these properties the equations

$$b_{\ell,m,i} = \frac{|S^{n-2}|}{C_m^\lambda(1)} d_{m,i}^\lambda \sum_{k=0}^{d+[m/2]} e_{k,\ell,m} a_{i,k,m},$$

where $0 \leq i, k \leq d + [m/2]$, determine uniquely $e_{k,\ell,m}$. This completes the proof. ■

4. Even dimensions

The nature of the problem changes considerable in even dimensions. On one hand, we do not have the easiness to handle polynomials, but on the other hand, just this sole inconvenience gives uniqueness on these spaces.

We have

$$(4.1) \quad (R_L f)_{l,m}(p) = |S^{n-2}| \int_1^\infty f_{\ell,m}(q) q^{2\lambda} \frac{D_m^\lambda(p/q)}{\sqrt{1-p^2/q^2}} dq,$$

where the dimension $n = 2\lambda + 2$ is even, $\lambda \geq 0$, $|S^0| = 2$ and

$$(4.2) \quad D_m^\lambda(x) = \begin{cases} \frac{C_m^\lambda(x)}{C_m^\lambda(1)} (1-x^2)^\lambda & \text{if } \lambda \geq 1 \\ \cos(m \arccos x) & \text{if } \lambda = 0 \end{cases}$$

is a polynomial of degree $m + 2\lambda$.

Lemma 4.1. *In the Taylor expansion*

$$\frac{D_m^\lambda(x)}{\sqrt{1-x^2}} = \sum_{i=0}^{\infty} d_{m,i}^\lambda x^{2i+m-2[m/2]},$$

the coefficients $d_{m,i}^\lambda$ are not zero. The series is convergent absolutely on $(-1, 1)$ and uniformly on $[-1 + \varepsilon, 1 - \varepsilon]$ for any $1 > \varepsilon > 0$.

Proof. For $\lambda > 0$, (8.932.1), (9.136.1) and (9.131.1) of [2] give

$$\frac{C_m^\lambda(x)}{(1-x^2)^{1/2-\lambda}} = K_{m,\lambda} F\left(\frac{1-m}{2}, \frac{1+m}{2}, \frac{1}{2}; x^2\right) + x K'_{m,\lambda} F\left(\frac{2-m}{2}, \frac{2+m}{2}, \frac{3}{2}; x^2\right),$$

where F is the Gauss hypergeometric function [2(9.1)], $K'_{m,\lambda} = 0$ for m even and $K_{m,\lambda} = 0$ for m odd.

For $\lambda = 0$ we prove that

$$\frac{\cos(m \arccos x)}{\sqrt{1-x^2}} = K_m F\left(\frac{1-m}{2}, \frac{1+m}{2}, \frac{1}{2}; x^2\right) + x K'_m F\left(\frac{2-m}{2}, \frac{2+m}{2}, \frac{3}{2}; x^2\right)$$

by means of (8.942.1), (9.136.1) and (9.131.1) from [2], where $K'_m = 0$ for m even and $K_m = 0$ for m odd.

The well known properties of the hypergeometric function and the above formulas prove the statement. ■

Let $b_{m,i}^\lambda = |S^{n-2}|d_{m,i}^\lambda$. According to Lemma 4.1 the Taylor expansion of (4.1) is

$$(4.3) \quad (R_L f)_{l,m}(p) = \sum_{i=0}^{\infty} b_{m,i}^\lambda p^{2i+m-2[m/2]} \int_1^{\infty} f_{\ell,m}(q) q^{2\lambda-2i+2[m/2]-m} dq.$$

Let $2\phi_{\ell,m}(x^2) = f_{\ell,m}(1/x)/x^n$. Changing the variable $q = 1/\sqrt{x}$ the integral in (4.3) becomes

$$(4.4) \quad c_{\ell,m,i} = \int_0^1 \phi_{\ell,m}(x) x^i x^{\varepsilon_m} dx.$$

Theorem 4.2. *The limited domain Radon transform R_L is injective on $L_{\alpha,\beta}^2(\mathbb{I}^n)$ if $n-1 \leq \alpha < n$ and $-1 < \beta < 0$.*

Proof. Easy calculation shows that $\phi_{\ell,m} \in L_{-1/2,0}^2([0,1])$. Then $R_L f(\omega, p) = 0$ for $p < 1$ implies $c_{\ell,m,i} = 0$ for all $i, m \geq 0$, hence $\phi_{\ell,m}$ must be zero. ■

Because (4.3) is an infinite series, we need more sophisticated tools to determine the range of R_L .

Theorem 4.3. *R_L maps the $L_{n-1,0}^2(\mathbb{I}^n)$ -closure of the span of the functions*

$$f_{k,\ell,m}(\omega, q) = q^{-n-2k} Y_{\ell,m}(\omega), \quad 0 \leq k$$

onto the $L_{\gamma,\delta}^2(\mathbb{I}^n)$ -closure of the span of the functions

$$F_{j,\ell,m}(\omega, p) = p^{2j+m-2[m/2]} Y_{\ell,m}(\omega) \quad 0 \leq j$$

continuously and bijectively.

Proof. The verification that $f_{k,\ell,m} \in L_{n-1,0}^2(\mathbb{I}^n)$ and $F_{j,\ell,m} \in L_{\gamma,\delta}^2(\mathbb{I}^n)$ is left to the reader. Since $R_L: L_{n-1,0}^2(\mathbb{I}^n) \rightarrow L_{\gamma,\delta}^2(\mathbb{I}^n)$ is continuous by Theorem 2.1, it maps a closed set into a closed set. Further, it is injective by Theorem 4.2, therefore we only have to give coefficients $e'_{k,\ell,m} \in \mathbb{R}$ so that

$$F_{j,\ell,m} = R_L(f_{\ell,m} Y_{\ell,m}), \quad \text{where} \quad f_{\ell,m} Y_{\ell,m} = \sum_{k=0}^{\infty} e'_{k,\ell,m} f_{k,\ell,m}.$$

The functions $f_{\ell,m}$ of this form can be written in the form

$$(4.5) \quad f_{\ell,m}(q) = \sum_{k=0}^{\infty} e_{k,\ell,m} q^{-n} P_k^{(0,\varepsilon_m)}(2q^{-2} - 1)$$

with unique coefficients $e_{k,\ell,m}$. To find these coefficients we consider (4.5) and (4.3), where the left hand side is substituted by $F_{j,\ell,m}$, as a system of linear equations with infinite dimensional matrix of entries

$$\begin{aligned} a_{i,k,m} &= \int_1^{\infty} q^{-2-2i+2[m/2]-m} P_k^{(0,\varepsilon_m)}(2q^{-2} - 1) dq \\ &= \frac{1}{2} \int_0^1 x^i P_k^{(0,\varepsilon_m)}(2x - 1) x^{\varepsilon_m} dx. \end{aligned}$$

By the orthogonality of $P_k^{(0,\varepsilon_m)}$ these are zero for $k > i$. For $i \geq k$

$$a_{i,k,m} = (k!)^2 \binom{i}{k} \prod_{h=0}^k \frac{1}{h + i + 1 + \varepsilon_m}$$

by [2(7.391.3)]. In virtue of the Stirling formula for the factorials we find the (not sharp) inequality

$$(4.6) \quad 16 < \lim_{i \rightarrow \infty} \frac{a_{i,i,m}}{\left(\frac{i}{4e}\right)^i}.$$

The system of equations

$$(4.7) \quad \delta_{i,j} = \begin{cases} 0 & \text{if } i \neq j \\ 1 & \text{if } i = j \end{cases} = b_{m,i}^{\lambda} \sum_{k=0}^i e_{k,\ell,m} a_{i,k,m}, \quad i, j \geq 0$$

obtained from (4.3), by substituting $F_{j,\ell,m}$ into the left hand side and substituting $f_{\ell,m}$ into the right hand side according to (4.5), determine the coefficients $e_{k,\ell,m}$.

What remains to prove is only that the function defined by (4.5) with these coefficients is in $L_{n-1,0}^2(\mathbb{E}^n)$, which is equivalent to

$$(4.8) \quad \int_1^{\infty} \left(\sum_{k=0}^{\infty} e_{k,\ell,m} P_k^{(0,\varepsilon_m)}(2q^{-2} - 1) \right)^2 q^{-2} dq < \infty.$$

Using Cauchy's inequality we can estimate the left hand side by

$$(4.9) \quad \sum_{k=0}^{\infty} (k+1)^2 e_{k,\ell,m}^2 \sum_{k=0}^{\infty} \int_1^{\infty} \left(\frac{P_k^{(0,\varepsilon_m)}(2q^{-2} - 1)}{k+1} \right)^2 q^{-2} dq.$$

First we estimate the second multiplier in (4.9). With a change in the variable we get

$$(4.10) \quad \int_1^\infty (P_k^{(0,\varepsilon m)}(2q^{-2} - 1))^2 q^{-2} dq = \frac{\sqrt{2}}{4} \int_{-1}^1 (P_k^{(0,\varepsilon m)}(x))^2 (1+x)^{-1/2} dx.$$

For m even, the right hand side evaluates explicitly $1/(4k+1)$ by [2(7.391.1)]. For m odd, $P_k^{(0,\varepsilon m)}(x) = C_k^{1/2}(x)$ by [2(8.962.4)], which is less than $C_k^{1/2}(1) = 1$. Therefore, the second multiplier in (4.9) is less than $\sum_{k=0}^\infty (k+1)^{-2} = \pi^2/6$.

To estimate the first multiplier in (4.9), we introduce for $N > j$ the N -dimensional vectors

$$\begin{aligned} \delta^N &= (b_{m,j}^\lambda)^{-1}(\delta_{0,j}, \dots, \delta_{i,j}, \dots, \delta_{N,j}) \\ e^N &= (e_{0,\ell,m}, 2e_{1,\ell,m}, \dots, (k+1)e_{k,\ell,m}, \dots, (N+1)e_{N,\ell,m}) \end{aligned}$$

and the $N \times N$ quadratic matrix $A^N = (a_{i,k,m}^{i,k,m})_{i,k=0}^N$. With these notations, (4.7) can be rewritten as $\delta^N = A^N e^N$ and therefore

$$(4.11) \quad \sum_{k=0}^N (k+1)^2 e_{k,\ell,m}^2 = \|e^N\|_2^2 \leq \|\delta^N\|_2^2 \cdot \|(A^N)^{-1}\|_2^2.$$

Obviously $\|\delta^N\|_2^2 = (b_{m,j}^\lambda)^{-2}$ does not depend on N . Since the matrix A^N is triangular so is its inverse $(A^N)^{-1}$, and the diagonal elements in $(A^N)^{-1}$ are $(i+1)/a_{i,i,m}$, which are also the eigenvalues of $(A^N)^{-1}$. In virtue of (4.6) the estimate

$$\|(A^N)^{-1}\|_2^2 \leq \max_{0 \leq i \leq M} \frac{i+1}{a_{i,i,m}}$$

is valid for M big enough. Together with (4.11) this completes the proof. ■

As a consequence we obtain the following.

Theorem 4.4. *The range of R_L is dense in $L_{\gamma,\delta}^2(\mathbb{B}^n)$, where $-1 < \gamma \leq 0$ and*

$$0 > \delta > \begin{cases} -1 & \text{if } \lambda > 0 \\ -1/2 & \text{if } \lambda = 0 \end{cases}$$

Proof. According to Theorem 4.3 it is enough to prove that $\langle F, F_{j,\ell,m} \rangle_{\gamma,\delta} = 0$ for all $j, \ell, m \geq 0$ implies $F = 0$ for $F \in L_{\gamma,\delta}^2(\mathbb{B}^n)$.

Let $\Phi(\omega, p^2) = F(\omega, p)$. Then $F \in L_{\gamma,\delta}^2(\mathbb{B}^n)$ implies $\Phi \in L_{(\gamma-1)/2,\delta}^2(\mathbb{B}^n)$, hence

$$0 = \langle F, F_{j,\ell,m} \rangle_{\gamma,\delta} = \int_0^1 \Phi_{\ell,m}(p) p^j p^{\frac{\gamma-1}{2} + \frac{m}{2} - [\frac{m}{2}]} (1-p)^\delta dp$$

for all $j \geq 0$ implies $\Phi_{\ell,m} \equiv 0$ for all $\ell, m \geq 0$, i.e. $F = 0$. ■

Note that we used the special values $\alpha = n - 1$ and $\beta = 0$ in Theorem 4.3 only to simplify the estimate of (4.10) that would need tedious calculations in general. However, the results can be extended to the spaces $L^2_{\alpha,\beta}$ for $n - 1 \leq \alpha < n$ and $-1 < \beta \leq 0$.

References

- [1] C. A. Berenstein, D. Walnut, Local inversion of the Radon transform in even dimensions using wavelets. *Proc. of the Conference 75 years of Radon transform*, International Press, 1994, 45–69.
- [2] I. S. Gradshteyn, I. M. Ryzhik, *Table of Integrals, Series and Products*. Acad. Press, New York, 1980.
- [3] S. Helgason, *The Radon Transform*. Birkhäuser, Boston, 1980.
- [4] W. R. Madych, Summability and approximate reconstruction from Radon transform data. *Cont. Math.* **113** (1990), 189–219.
- [5] E. T. Quinto, The invertibility of rotation invariant transforms. *J. Math. Anal. Appl.* **91** (1983), 510–522.
- [6] E. T. Quinto, Singular value decompositions and inversion methods for the exterior Radon transform and a spherical transform. *J. Math. Anal. Appl.* **95** (1983), 437–448.
- [7] R. Seeley, Spherical harmonics. *Amer. Math. Monthly* **73** (1966), 115–121.

Bolyai Institute, Szeged University
 Aradi vértanúk tere 1.
 H-6720 Szeged, HUNGARY
 Kurusa@math.u-szeged.hu

Received 20.01.1996