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Extensions of a Symmetric Second-Order Differential Operator

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Presented by Bl. Sendov

Let L_0 be the closure of the expression of the differential operator

$$l(y) = -y''(x) + \frac{\nu^2 - \frac{1}{4}}{x^2} y(x) + q(x) y(x), \quad 0 < x < \infty.$$

where $\nu \geq 1$ and $q(x)$ is a real continuous function in $[0, \infty)$. In this case, the defect index of the operator L_0 is $(0, 0)$ or $(1, 1)$. For the case of defect index $(1, 1)$, the space of boundary values of the operator L_0 and all of its maximal dissipative, maximal accretive and selfadjoint extensions are obtained.

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1. Introduction

Let us consider the following differential operator

$$(1) \quad l(y) = -y''(x) + \frac{\nu^2 - \frac{1}{4}}{x^2} y(x) + q(x) y(x), \quad 0 < x < \infty,$$

where $\nu \geq 1$ and $q(x)$ is a real continuous function in $[0, \infty)$.

Let us denote by L_0 the closure of the minimal symmetric operator ([1]) generated by (1). We denote by D the set of all the functions $y(x)$ from $L_2(0, \infty)$ whose first derivatives are locally absolutely continuous in $(0, \infty)$ and $l(y) \in L_2(0, \infty)$; D is the domain of the maximal operator L , and $L = L_0^*$, [1]. The symmetric operator L_0 has defect index $(0, 0)$ or $(1, 1)$, [1]. For defect index $(0, 0)$ the operator L_0 is selfadjoint, that is, $L_0^* = L_0 = L$.

Previously, for $\nu = \frac{1}{2}$, the space of the boundary values of operator L_0 , maximal dissipative and selfadjoint extensions of the operator are studied in [2].

In this work, we assume that the symmetric operator L_0 has defect index (1, 1).

Let $v_1(x)$ and $v_2(x)$ denote the solutions of $l(y) = 0$ satisfying the initial conditions

$$v_1(0) = 1, \quad v_1'(0) = 1, \quad v_2(0) = 1, \quad v_2'(0) = 1.$$

Clearly, $v_1(x)$ and $v_2(x)$ are linearly independent and their Wronskians are equal to one:

$$W[v_1, v_2]_x = W[v_1, v_2]_1 = 1, \quad (1 \leq x \leq \infty).$$

We recall that a triple $(\mathcal{H}, \Gamma_1, \Gamma_2)$, where \mathcal{H} is a Hilbert space and Γ_1 and Γ_2 are linear maps of $D(A^*)$ into \mathcal{H} , is called the space of boundary values (SBV) of closed symmetric operator A in the Hilbert space \mathcal{H} with equal defect indices, if the following two conditions hold:

1) For every $f, g \in D(A^*)$

$$(A^* f, g)_H - (f, A^* g)_H = (\Gamma_1 f, \Gamma_2 g)_H - (\Gamma_2 f, \Gamma_1 g)_H,$$

2) For every $g_1, g_2 \in \mathcal{H}$ there is a vector $f \in D(A^*)$.

We consider the following linear maps of D into \mathbb{C} :

$$(2) \quad \Gamma_1 f = W[f, v_1]_\infty, \quad \Gamma_2 f = W[f, v_2]_\infty, \quad f \in D.$$

2. Results

Theorem 1. *The triple $(\mathbb{C}, \Gamma_1, \Gamma_2)$ defined by (2) is the space of boundary values of the operator L_0 .*

Proof. In order to check the first condition of (SBV), we first prove the following lemma.

Lemma 1. *For arbitrary functions $y(x), z(x) \in D$ we have*

$$(3) \quad W[y, \bar{z}]_x = W[y, v_1]_x \cdot W[\bar{z}, v_2]_x - W[y, v_2]_x \cdot W[\bar{z}, v_1]_x, \quad 0 \leq x \leq \infty.$$

Proof. Observing that $v_1(x)$ and $v_2(x)$ are real functions, we have,

$$\begin{aligned} & W[y, v_1]_x \cdot W[\bar{z}, v_2]_x - W[y, v_2]_x \cdot W[\bar{z}, v_1]_x \\ &= (y(x)v_1'(x) - y'(x)v_1(x))(\bar{z}(x)v_2'(x) - \bar{z}'(x)v_2(x)) \\ &\quad - (y(x)v_2'(x) - y'(x)v_2(x))(\bar{z}(x)v_1'(x) - \bar{z}'(x)v_1(x)) \\ &= y(x)\bar{z}'(x) - y'(x)\bar{z}(x) = W[y, \bar{z}]_x, \quad 0 \leq x \leq \infty. \end{aligned}$$

Thus, Lemma 1 is proved. \blacksquare

For $\nu \geq 1$ and operator L_0 having defect index $(1, 1)$, and for every $y(x), z(x) \in D$ satisfying $W[y, \bar{z}]_0 = 0$, the following Lagrange formula is satisfied:

$$(4) \quad (Ly, z)_{L_2(0, \infty)} - (y, Lz)_{L_2(0, \infty)} = W[y, z]_\infty$$

Using Lemma 1, from (4) the following relation is obtained

$$\begin{aligned} (\Gamma_1 y, \Gamma_2 z)_{\mathbb{C}} - (\Gamma_2 y, \Gamma_1 z)_{\mathbb{C}} &= W[y, v_1]_\infty \cdot W[\bar{z}, v_2]_\infty - W[y, v_2]_\infty \cdot W[\bar{z}, v_1]_\infty \\ &= W[y, \bar{z}]_\infty. \end{aligned}$$

Therefore, the first requirement of the SBV is fulfilled. The second requirement is proved by the use of the following lemma.

Lemma 2. For any complex numbers β_0, β_1 , there is a function $y \in D$ satisfying

$$(5) \quad W[y, v_1]_\infty = \beta_0, \quad W[y, v_2]_\infty = \beta_1.$$

Proof. Let us denote by L_1 the closure of the minimal symmetric operator generated by $l(y)$ in $1 \leq x < \infty$. For any complex numbers $\gamma_0, \gamma_1, \beta_0$ and β_1 , there is a function $y_1(x) \in D(L_1^*)$ which satisfies the following conditions

$$(6) \quad y(1) = \gamma_0, \quad y'(1) = \gamma_1, \quad W[y, v_1]_\infty = \beta_0, \quad W[y, v_2]_\infty = \beta_1.$$

Now, let us prove these relations. We consider a function $f(x) \in L_2(1, \infty)$ satisfying

$$(7) \quad (f, v_1)_{L_2(1, \infty)} = \beta_0 + \gamma_1, \quad (f, v_2)_{L_2(1, \infty)} = \beta_1 - \gamma_0.$$

Let $y_1(x)$ denote the solution of equation $l(y) = f(x)$, $(1 \leq x < \infty)$ satisfying the initial conditions

$$y(1) = \gamma_0, \quad y'(1) = \gamma_1.$$

This solution can be written as

$$y_1(x) = \gamma_0 v_1(x) + \gamma_1 v_2(x) + \int_1^x \{v_1(x) v_2(\xi) - v_1(\xi) v_2(x)\} f(\xi) d\xi.$$

This expression shows that $y_1(x) \in D(L_1^*)$. Let us apply Lagrange's formula to the functions $y_1(x)$ and $v_j(x)$ ($j = 1, 2$)

$$(8) \quad (f, v_j)_{L_2(1, \infty)} = (l(y_1), v_j)_{L_2(1, \infty)} = W[y_1, v_j]_\infty - W[y_1, v_j]_1 + (y, l(v_j))_{L_2(1, \infty)}.$$

If we let

$$l(v_j) = 0, \quad y_1(1) = \gamma_0, \quad y_1'(1) = \gamma_1$$

in (8), we find

$$W[y_1, v_j] = \begin{cases} -\gamma_1 & \text{for } j = 1 \\ \gamma_0 & \text{for } j = 2 \end{cases}$$

and

$$(f, v_1)_{L_2(1, \infty)} = W[y_1, v_1]_{\infty} + \gamma_1,$$

$$(f, v_2)_{L_2(1, \infty)} = W[y_1, v_2] - \gamma_0.$$

From (7) we obtain

$$W[y_1, v_1]_{\infty} = \beta_0, \quad W[y_1, v_2]_{\infty} = \beta_1$$

Hence, we have proved that there exists a function $y_1 \in D(L_1^*)$ which satisfies (6). For any complex numbers γ_0 and γ_1 , let

$$y_1(x) = \gamma_0 v_1(x) + \gamma_1 v_2(x), \quad (0 < x \leq 1).$$

Then, let us define

$$y(x) = \begin{cases} y_2(x) & 0 < x \leq 1 \\ y_1(x) & 1 \leq x < \infty. \end{cases}$$

It is clear that $y \in D$. With respect to the condition (6) we obtain

$$W[y, v_1]_{\infty} = W[y_1, v_1]_{\infty} = \beta_0, \quad W[y, v_2]_{\infty} = W[y_1, v_2]_{\infty} = \beta_1.$$

Hence, Lemma 2 and Theorem 1 are proved. Using Theorem 1 and Theorem 1.6 [3], we can state the following theorem.

Theorem 2. For every number $h \in \mathbb{C}$, $\text{Im } h \geq 0$ or $h = \infty$, the restriction of L to the set of functions $y \in D$ satisfying either

$$(9) \quad W[y, v_1]_{\infty} - h W[y, v_2]_{\infty} = 0,$$

or

$$(10) \quad W[y, v_1]_{\infty} + h W[y, v_2]_{\infty} = 0,$$

is respectively the maximal dissipative and accretive extension of the symmetric operator L_0 . Conversely, every maximally dissipative (accretive) extension of L_0 is the compression (restriction) of L to the set of functions $y(x) \in D$ satisfying (9), (10). If $\text{Im } h = 0$ or $h = \infty$, the conditions (9) and (10) define

selfadjoint extensions of the operator L_0 . For $h = \infty$, the conditions (9), (10) form $W[y, v_2]_\infty = 0$.

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