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ON THE CLASSIFICATION OF UNITARY LOWEST WEIGHT REPRESENTATIONS OF CONFORMAL CURRENT ALGEBRAS

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*To the memory of my teacher
Yaroslav Tagamlitzki*

1. Introduction. In searching for the extreme elements of a partially ordered set K the good half of the work consists in devising a decomposition of a generic positive element of K into a sum (or a convex linear combination) of positive elements [1]. The purpose of this note is to point out the use of such a decomposition in a context, in which K is not a convex set.

The problem is concerned with the classification of the unitary lowest weight representations of an infinite dimensional Lie algebra VG , which appears in conformal quantum field theory models on the circle. In the rest of this introduction I shall define the algebra and say some words about its physical interpretation (for more detail see, e. g. [2]).

Let G be a simple compact Lie group and let d_G be its Lie algebra. To each element $X \in d_G$ we make correspond a sequence of elements $X_n, n \in \mathbb{Z}$ of an infinite dimensional "affine Kac-Moody Lie algebra" \widehat{d}_G in such a way that the following commutation relations hold:

$$[X_m, Y_n] = [X, Y]_{n+m} + kn(X, Y)\delta_{n,-m}$$

where (X, Y) is a suitably normalized invariant inner product (the positive "Killing form") on d_G , and k is a "central charge"

$$[k, X_n] = 0.$$

Thus, \widehat{d}_G appears as a \mathbb{Z} -graded infinite Lie algebra which contains d_G as a subalgebra (if we identify X with X_0).

The Virasoro algebra Vir is a central extension of the algebra of first-order differential operators that generate the diffeomorphisms of the circle. It is another \mathbb{Z} -graded algebra with generators L_n and c satisfying

$$(1.3) \quad [L_n, L_m] = (n-m)L_{n+m} + \frac{c}{12} n(n^2-1)\delta_{n,-m}$$

$$(1.4) \quad [c, L_n] = 0.$$

The conformal current algebra VG is defined as the semidirect sum of Vir with \widehat{d}_G , in which the commutation relations (1.1-1.4) are supplemented by

$$(1.5) \quad [X_m, L_n] = mX_{n+m}.$$

The physical context in which such an algebra appears is an 1-dimensional conformal quantum field theory on the circle with an internal symmetry group G . The fields of such a theory are defined in general as *op-*
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rator-valued distributions on a Hilbert space \mathcal{H} (satisfying certain requirements). In our case they include the stress energy tensor

$$(1.6) \quad T(z) = \sum_n L_n z^{-n-2}, \quad L_n^* = L_{-n} (n \in \mathbb{Z})$$

and the current

$$(1.7) \quad X(z) = \sum_n X_n z^{-n-1}, \quad X_n^* = X_{-n}$$

(X belongs to a d_G dimensional space, where $d_G = \dim G$). The states of the theory include the vacuum vector $|0\rangle \in \mathcal{H}$ satisfying $\langle 0|0\rangle = 1$ and

$$(1.8) \quad L_n |0\rangle = 0 \text{ for } n \geq -1, \quad X_m |0\rangle = 0 \text{ for } m \geq 0.$$

We make further the following irreducibility assumption: if a bounded operator B in \mathcal{H} commutes with all the fields of the theory, then it is a multiple of the identity in \mathcal{H} . In particular, the central charges c and k act as multiplication of real numbers.

Corollaries. 1. The positivity of the inner product in \mathcal{H} and the hermiticity properties in (1.6) and (1.7) imply that c and k are positive

$$(1.9) \quad 0 < \|L_{-2}|0\rangle\|^2 = \langle 0|L_2 L_{-2}|0\rangle = \frac{c}{2}$$

$$(1.10) \quad 0 < \|X_{-1}|0\rangle\|^2 = \langle 0|X_1 X_{-1}|0\rangle = k(X, X).$$

2. Eqs. (1.6-1.8) allow to find all correlation functions of the stress-energy tensor and the current:

$$(1.11) \quad \langle 0|T(z_1)T(z_2)|0\rangle = \frac{c}{2} z_{12}^{-4}, \quad z_{12} = z_1 - z_2.$$

$$(1.12) \quad \langle 0|X(z_1)Y(z_2)|0\rangle = \frac{k}{z_{12}^2} (X, Y)$$

$$(1.13) \quad \langle 0|X(z_1)Y(z_2)T(z_3)|0\rangle = k(X, Y) z_{13}^{-2} z_{23}^{-2} \text{ etc.}$$

The operator L_0 is interpreted physically as the energy operator. Energy positivity singles out lowest weight (LW) representations of VG. Any LW unitary irreducible representation (UIR) of VG is characterized by a LW vector $|\Psi\rangle$ such that

$$(1.14) \quad L_n |\Psi\rangle = 0 = X_n |\Psi\rangle \text{ for } n = 1, 2, \dots, (L_0 - \Delta) |\Psi\rangle > 0$$

(which is usually also assumed to be a highest weight eigenvector of a Cartan subalgebra of d_G).

The classification of LWUIRs of VG has only been completed recently [3], following developments in [4-7] among others. We shall review in Sec. 2 a central point of this derivation, related to a decomposition of $T(z)$.

2. LWUIRs of VG. Proposition 1. Given any unitary LW representation of d_G , it can be extended to a LW representation of VG by setting

$$(2.1) \quad L_{2n}^G = \frac{1}{k+C_2} \left\{ \vec{Q}_n^2 + 2 \sum_{n=1}^{\infty} \vec{Q}_{n-m} \vec{Q}_{n+m} \right\}, \quad L_{2n+1}^G = \frac{2}{k+C_2} \sum_{l=0}^{\infty} \vec{Q}_{n-m} \vec{Q}_{n+1+m},$$

where $\vec{Q} = (Q_1, \dots, Q^{d_G})$ is an orthonormal basis of d_G so that

$$(2.2) \quad (Q^a, Q^b) = \frac{1}{2} \delta_{ab}, \quad [Q_n^a, Q_m^b] = if_{abc} Q_{n+m}^c + \frac{kn}{2} \delta_{n-m} \delta_{ab}$$

and G_2 is the eigenvalue of the second order Casimir operator for the adjoint representation of d_G

$$(2.3) \quad f_{sta}f_{stb} = C_2\delta_{ab}.$$

The Virasoro central charge associated with L_n^G is

$$(2.4) \quad c_G (= 2[L_2^G, L_{-2}^G] - 8L_0^G) = \frac{k}{k+C_2} d_G.$$

The proof (2.6) involves a straightforward (albeit somewhat lengthy) computation using (2.1-2.3) (One should first verify $[Q_m^a, L_n^G] = m Q_{m+n}^a$ and then (1.3), (2.4)).

Given a LWUIR of d_G , the construction of Proposition 1 gives a minimal representation of VG in the sense of the following result.

Proposition 2. Given an arbitrary LWUIR of VG the operators

$$(2.5) \quad l_n = L_n - L_n^G$$

give rise to a LWUIR of Vir with central charge

$$(2.6) \quad c_l = c - c_G \geq 0.$$

Proof. It follows from (1.5) and from Proposition 1 that

$$(2.7) \quad [Q_m^a, l_n] = 0 = [L_m^G, l_n].$$

The rest is straightforward. The inequality (2.6) is a consequence of the assumed unitarity of the representation under consideration.

Propositions 1 and 2 reduce the study of the LWUIRs of VG to the known classification of such representations for d_G and Vir (2.4-2.7).

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