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CANONICAL CONNECTION AND THE CANONICAL CONFORMAL GROUP ON A RIEMANNIAN ALMOST-PRODUCT MANIFOLD

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On an almost Hermitian manifold (M, J, g) there exists a unique linear connection D with torsion tensor T such that $D_X J = D_X g = 0$ and $T(x, JY) = T(JX, Y)$ for all vector fields X, Y on M . This is the Hermitian connection on the manifold [1, 2]. Another proof of this result has been given in [3]. The group of the conformal transformations of the metric g generates the conformal group of transformations of D .

On an almost complex manifold with B -metric (M, J, g) there exists a unique canonical connection D such that $D_X J = D_X g = 0$, or equivalently $D_X g = 0, D_X \tilde{g} = 0$ for any vector field X on M [4]. Here \tilde{g} is the associated metric to the metric g . Considering the general group of the conformal transformations of the B -metric g , in [4] there is also obtained the canonical conformal group of transformations of the canonical connection D and its remarkable subgroups.

In this paper we treat analogous problems on Riemannian almost-product manifolds.

1. Preliminaries. Let (M, g) be a differentiable manifold with metric g .

Definition 1. An almost-product structure on M is a $(1,1)$ -tensor field v of constant rank for which $v^2 = v$.

If we denote $h = I - v$ and $\mathcal{P} = v - h$, then $h^2 = h$, so that h is also an almost-product structure on M and $\mathcal{P}^2 = I$. It is easy to be proved that

$$(1) \quad vh = hv = 0, \quad v = \frac{1}{2}(I + \mathcal{P}), \quad h = \frac{1}{2}(I - \mathcal{P}).$$

Definition 2. A Riemannian almost-product structure on M is an almost-product structure on M such that

$$(2) \quad g(X, Y) = g(vX, vY) + g(hX, hY),$$

where X, Y are arbitrary vector fields on M .

From (1) and (2) it follows immediately that

$$(3) \quad g(vX, Y) = g(X, vY), \quad g(hX, Y) = g(X, hY), \quad g(vX, hY) = 0,$$

and

$$(4) \quad g(\mathcal{P}X, \mathcal{P}Y) = g(X, Y)$$

for all differentiable vector fields X, Y on M .

In this case we shall say that \mathcal{P} and g are compatible, and that \mathcal{P} is the tensor field defining a Riemannian almost-product structure on a Riemannian manifold M with metric tensor field g .

Because of (4) and (3) we can also define a symmetric 2-covariant tensor field \tilde{g} on M by

$$(5) \quad \tilde{g}(X, Y) = g(\mathcal{P}X, Y).$$

In what follows (M, \mathcal{P}, g) will be a $2n$ -dimensional Riemannian almost-product manifold, i. e. \mathcal{P} will be the tensor field defining the Riemannian almost-product structure on M , and g will be the metric tensor field on M such that $\mathcal{P}^2 = I$, $g(\mathcal{P}X, Y) = g(X, \mathcal{P}Y)$ for all vector fields X, Y on M . The associated metric tensor field \tilde{g} on the manifold is given by (5). The metric \tilde{g} is necessarily of signature (n, n) .

Further X, Y, Z will stand for arbitrary differentiable vector fields on M .

The Levi-Civita connections of g and \tilde{g} will be denoted by ∇ and $\tilde{\nabla}$, respectively.

The difference $\tilde{\nabla}_X Y - \nabla_X Y$ is a tensor field of type $(1, 2)$ on M and it will be denoted by

$$(6) \quad \Phi(X, Y) = \tilde{\nabla}_X Y - \nabla_X Y.$$

This is the fundamental tensor field of the manifold. The corresponding tensor field of type $(0, 3)$ will be denoted by the same letter: $\Phi(X, Y, Z) = g(\Phi(X, Y), Z)$. The fundamental tensor field has the following symmetries:

$$(7) \quad \begin{aligned} \Phi(X, Y) &= \Phi(Y, X) \\ \Phi(X, Y, Z) + \Phi(X, \mathcal{P}Y, \mathcal{P}Z) + \Phi(Z, X, Y) + \Phi(\mathcal{P}Z, X, \mathcal{P}Y) &= 0, \\ \Phi(X, Y, Z) + \Phi(\mathcal{P}X, \mathcal{P}Y, Z) + \Phi(\mathcal{P}X, Y, \mathcal{P}Z) + \Phi(X, \mathcal{P}Y, \mathcal{P}Z) &= 0. \end{aligned}$$

2. Partial decomposition of the space of torsion tensors. The canonical linear connection in our considerations will have a torsion. Thus, we have to study the properties of the torsion tensors.

Let $(V, \mathcal{P}, g) = T_p M$, $p \in M$ and \mathcal{F} be the vector space of all tensors T of type $(0, 3)$ over V with the property $T(x, y, z) = -T(y, x, z)$, $x, y, z \in V$. The metric g induces on \mathcal{F} an inner product $\langle \cdot, \cdot \rangle$ in the following way:

$$\langle T', T'' \rangle = g^{i^q} g^{j^r} g^{k^s} T'(l_i, l_j, l_k) T''(l_q, l_r, l_s), \quad T', T'' \in \mathcal{F},$$

where $\{l_i\}$ ($i = 1, \dots, 2n$) is a basis of V .

The natural representation of the group $G = \{a \in O(2n) | a = \begin{pmatrix} A & B \\ B & A \end{pmatrix}, A, B \text{ are } n \times n\text{-matrices}\}$ in V induces a representation λ of G in \mathcal{F} :

$$(\lambda a)T(x, y, z) = T(a^{-1}x, a^{-1}y, a^{-1}z), \quad a \in G, T \in \mathcal{F}, x, y, z \in V,$$

so that $\langle (\lambda a)T', (\lambda a)T'' \rangle = \langle T', T'' \rangle$, $a \in G, T', T'' \in \mathcal{F}$.

Let L be the linear operator $L: \mathcal{F} \rightarrow \mathcal{F}$ defined by

$$L(T)(x, y, z) = T(\mathcal{P}x, \mathcal{P}y, z), \quad T \in \mathcal{F}, x, y, z \in V.$$

The following proposition follows by simple computations.

Lemma 1. *The operator L is an involutive isometry of \mathcal{F} and commutes with the action of G , i. e. $L^2 = id$,*

$$\langle L(T'), L(T'') \rangle = \langle T', T'' \rangle, \quad L((\lambda a)T) = (\lambda a)(L(T)),$$

where $T, T', T'' \in \mathcal{F}$, $a \in G$.

This lemma implies that L has two eigenvalues (± 1) and the corresponding eigenspaces

$$\mathcal{F}^- = \{T \in \mathcal{F} | L(T) = -T\}, \quad \mathcal{F}^+ = \{T \in \mathcal{F} | L(T) = T\}$$

are invariant orthogonal subspaces of \mathcal{F} .

In order to decompose \mathcal{F}^- we consider the linear operator $L_1: \mathcal{F}^- \rightarrow \mathcal{F}^-$ defined by

$$L_1(T)(x, y, z) = \frac{1}{2} \{T(y, z, x) + T(z, x, y) - T(\mathcal{P}y, z, \mathcal{P}x) - T(z, \mathcal{P}x, \mathcal{P}y)\},$$

where $T \in \mathcal{F}^-$, $x, y, z \in V$. We have

Lemma 2. *The operator L_1 is an involutive isometry and commutes with the action of G .*

This lemma implies that the eigenspaces

$$\mathcal{F}_1 = \{T \in \mathcal{F}^- \mid L_1(T) = -T\}, \quad \mathcal{F}_2 = \{T \in \mathcal{F}^- \mid L_1(T) = T\}$$

are invariant and orthogonal subspaces of \mathcal{F}^- .

To decompose \mathcal{F}^+ we define the linear operator $L_2: \mathcal{F}^+ \rightarrow \mathcal{F}^+$ in the following way

$$L_2(T)(x, y, z) = T(\mathcal{P}x, y, \mathcal{P}z), \quad T \in \mathcal{F}^+, \quad x, y, z \in V.$$

We have

Lemma 3. *The operator L_2 is an involutive isometry and commutes with the action of G .*

This lemma implies that the eigenspaces

$$\mathcal{F}_3 = \{T \in \mathcal{F}^+ \mid L_2(T) = -T\}, \quad \mathcal{F}_4 = \{T \in \mathcal{F}^+ \mid L_2(T) = T\}$$

are invariant and orthogonal subspaces of \mathcal{F}^+ .

From the definitions of the spaces \mathcal{F}_k ($k=1, 2, 3, 4$) we obtain

Lemma 4. *Let $T \in \mathcal{F}$. Then*

- 1) $T \in \mathcal{F}_1$ iff $T(\mathcal{P}x, \mathcal{P}y, z) = -T(x, y, z)$, $T(x, y, z) + T(y, z, x) + T(z, x, y) = 0$;
- 2) $T \in \mathcal{F}_2$ iff $T(\mathcal{P}x, \mathcal{P}y, z) = -T(x, y, z)$, $T(\mathcal{P}x, y, z) + T(\mathcal{P}y, z, x) + T(\mathcal{P}z, x, y) = 0$;
- 3) $T \in \mathcal{F}_3$ iff $T(\mathcal{P}x, \mathcal{P}y, z) = T(x, y, z) = -T(\mathcal{P}x, y, \mathcal{P}z)$, or equivalently $T(\mathcal{P}x, y, z) = T(x, \mathcal{P}y, z) = -T(x, y, \mathcal{P}z)$;
- 4) $T \in \mathcal{F}_4$ iff $T(\mathcal{P}x, \mathcal{P}y, z) = T(x, y, z) = T(\mathcal{P}x, y, \mathcal{P}z)$, or equivalently $T(\mathcal{P}x, y, z) = T(x, \mathcal{P}y, z) = T(x, y, \mathcal{P}z)$.

Using Lemmas 1, 2 and 3 we obtain the following decomposition of \mathcal{F} :

Theorem 1. $\mathcal{F} = \mathcal{F}_1 \oplus \mathcal{F}_2 \oplus \mathcal{F}_3 \oplus \mathcal{F}_4$, where \mathcal{F}_i ($i=1, 2, 3, 4$) are invariant orthogonal subspaces of \mathcal{F} .

The projection operators of \mathcal{F} in \mathcal{F}_i ($i=1, 2, 3, 4$) are given in the following theorem.

Theorem 2. *Let $T \in \mathcal{F}$ and denote by p_i ($i=1, 2, 3, 4$) the projection operators of \mathcal{F} in \mathcal{F}_i . Then*

$$p_1(x, y, z) = \frac{1}{8} \{2T(x, y, z) - T(y, z, x) - T(z, x, y) + T(\mathcal{P}y, z, \mathcal{P}x) + T(z, \mathcal{P}x, \mathcal{P}y) - 2T(\mathcal{P}x, \mathcal{P}y, z) + T(\mathcal{P}y, \mathcal{P}z, x) + T(\mathcal{P}z, \mathcal{P}x, y) - T(y, \mathcal{P}z, \mathcal{P}x) - T(\mathcal{P}z, x, \mathcal{P}y)\}$$

$$p_2(x, y, z) = \frac{1}{8} \{2T(x, y, z) + T(y, z, x) + T(z, x, y) - T(\mathcal{P}y, z, \mathcal{P}x) - T(z, \mathcal{P}x, \mathcal{P}y) - 2T(\mathcal{P}x, \mathcal{P}y, z) - T(\mathcal{P}y, \mathcal{P}z, x) - T(\mathcal{P}z, \mathcal{P}x, y) + T(y, \mathcal{P}z, \mathcal{P}x) + T(\mathcal{P}z, x, \mathcal{P}y)\},$$

$$p_3(x, y, z) = \frac{1}{4} \{T(x, y, z) + T(\mathcal{P}x, \mathcal{P}y, z) - T(\mathcal{P}x, y, \mathcal{P}z) - T(x, \mathcal{P}y, \mathcal{P}z)\},$$

$$p_4(x, y, z) = \frac{1}{4} \{T(x, y, z) + T(\mathcal{P}x, \mathcal{P}y, z) + T(\mathcal{P}x, y, \mathcal{P}z) + T(x, \mathcal{P}y, \mathcal{P}z)\}.$$

Proof. We will compute for example v_1 . The other projections can be obtained in a similar way.

From Lemma 1 it follows that the tensor $\frac{1}{2}(T-L(T))$ is the projection of T into $\mathcal{T} = \mathcal{T}_1 \oplus \mathcal{T}_2$. Using Lemma 2 we find

$$p_1 = \frac{1}{4} \{T - L(T) - L_1(T) + L_1(L(T))\},$$

which gives the required expression of p_1 .

We need the following characteristic of the class $\mathcal{T}_1 \oplus \mathcal{T}_3$.

Lemma 5. Let $T \in \mathcal{T}$. Then $T \in \mathcal{T}_1 \oplus \mathcal{T}_3$ iff

$$(8) \quad T(x, y, z) + T(y, z, x) + T(\mathcal{P}x, y, \mathcal{P}z) + T(y, \mathcal{P}z, \mathcal{P}x) = 0.$$

Proof. Let $T \in \mathcal{T}_1 \oplus \mathcal{T}_3$, i. e. $T = p_1 + p_3$. It can be directly seen that p_1 and p_3 satisfy (8).

Conversely, let T satisfy (8). Since $T = T' + T''$, where T' and T'' are the components of T in $\mathcal{T}_1 \oplus \mathcal{T}_2$ and $\mathcal{T}_3 \oplus \mathcal{T}_4$ respectively, the substitutions of T' and T'' into (8) imply $T' = p_1$ and $T'' = p_3$, i. e. $T = p_1 + p_3$.

3. Natural connection on a Riemannian almost-product manifold.

Definition 3. A linear connection D on the Riemannian almost-product manifold (M, \mathcal{P}, g) is said to be natural if $D\mathcal{P} = 0$ and $Dg = 0$.

Since for an arbitrary linear connection

$$(D_x \tilde{g})(Y, Z) = (D_x g)(\mathcal{P}Y, Z) + g((D_x \mathcal{P})Y, Z),$$

it follows immediately that the linear connection D is natural iff $Dg = D\tilde{g} = 0$.

In this section we shall study the set of the natural connections on a Riemannian almost-product manifold (M, \mathcal{P}, g) .

Let T be the torsion tensor of a natural connection D . We denote by the same letter the corresponding tensor of type $(0, 3)$; $T(X, Y, Z) = g(T(X, Y), Z)$. The corresponding tensors p_i ($i = 1, 2, 3, 4$) associated with T are defined in Section 2.

Theorem 3. The linear connection D with torsion tensor T is natural iff

$$(9) \quad v_1(X, Y, Z) = \frac{1}{4} \{-\Phi(X, Y, Z) + \Phi(Y, Z, X) - \Phi(X, \mathcal{P}Y, \mathcal{P}Z) - \Phi(Y, \mathcal{P}Z, \mathcal{P}X) + 2\Phi(Z, \mathcal{P}X, \mathcal{P}Y)\},$$

$$(10) \quad p_3(X, Y, Z) = -\frac{1}{2} \{\Phi(Z, X, Y) + \Phi(Z, \mathcal{P}X, \mathcal{P}Y)\}.$$

Proof. We denote

$$(11) \quad D_x Y - \nabla_x Y = S(X, Y).$$

Since $\nabla g = 0$, the condition $Dg = 0$ is equivalent to

$$g(S(X, Y), Z) + g(S(X, Z), Y) = 0.$$

Hence

$$(12) \quad g(S(X, Y), Z) = \frac{1}{2} \{T(X, Y, Z) - T(Y, Z, X) + T(Z, X, Y)\},$$

which is the Theorem of Hayden [5].

Further we denote

$$(13) \quad D_x Y - \tilde{\nabla}_x Y = \tilde{S}(X, Y).$$

Analogously, because of $\tilde{\nabla} g \approx 0$, the condition $Dg \approx 0$ is equivalent to

$$g(\tilde{S}(X, Y), \mathcal{P}Z) + g(\tilde{S}(X, Z), \mathcal{P}Y) = 0.$$

Then

$$(14) \quad g(\tilde{S}(X, Y), Z) = \frac{1}{2} \{T(X, Y, Z) - T(Y, \mathcal{P}Z, \mathcal{P}X) + T(\mathcal{P}Z, X, \mathcal{P}Y)\}.$$

Using (11), (13) and (6), we get

$$\Phi(X, Y, Z) = g(S(X, Y, Z)) - g(\tilde{S}(X, Y), Z).$$

From the last equality and (12), (14) it follows

$$(15) \quad \Phi(X, Y, Z) = \frac{1}{2} \{-T(Y, Z, X) + T(Z, X, Y) + T(Y, \mathcal{P}Z, \mathcal{P}X) - T(\mathcal{P}Z, X, \mathcal{P}Y)\}.$$

From Theorem 1 and (15), using Lemma 4, we get

$$\begin{aligned} \Phi(X, Y, Z) = \frac{1}{2} \{ & -p_1(Y, Z, X) + p_1(Z, X, Y) + p_1(Y, \mathcal{P}Z, \mathcal{P}X) - p_1(\mathcal{P}Z, X, \mathcal{P}Y) \\ & - 2p_3(Y, Z, X) + 2p_3(Z, X, Y) \}. \end{aligned}$$

From the last formula we easily find (9) and (10).

For the inverse, let the projections p_1 and p_3 of T be (9) and (10). Substituting this projections into $T = p_1 + p_2 + p_3 + p_4$ we get (15), which implies D is natural.

Theorem 4. Let $Q'(X, Y, Z)$ and $Q''(X, Y, Z)$ be tensor fields on (M, \mathcal{P}, g) having the properties:

- (i) $Q'(X, Y, Z) = -Q'(Y, X, Z) = -Q'(\mathcal{P}X, \mathcal{P}Y, Z),$
 $Q'(\mathcal{P}X, Y, Z) + Q'(\mathcal{P}Y, Z, X) + Q'(\mathcal{P}Z, X, Y) = 0,$
- (ii) $Q''(X, Y, Z) = -Q''(Y, X, Z), Q''(\mathcal{P}X, Y, Z) = Q''(X, \mathcal{P}Y, Z) = Q''(X, Y, \mathcal{P}Z).$

Then there exists a unique natural connection D such that the components p_2 and p_4 of its torsion tensor T are the given tensors Q' and Q'' , respectively.

Proof. Existence. Taking into account Theorem 3, we construct the connection D

$$\begin{aligned} g(D_X Y, Z) = g(-X, Y, Z) + \frac{1}{4} \{ & \Phi(X, Y, Z) - 2\Phi(Z, X, Y) - \Phi(X, \mathcal{P}Y, \mathcal{P}Z) \} \\ & + \frac{1}{2} \{ Q'(X, Y, Z) - Q'(Y, Z, X) + Q'(Z, X, Y) \} \\ & + \frac{1}{2} \{ Q''(X, Y, Z) - Q''(Y, Z, X) + Q''(Z, X, Y) \}. \end{aligned}$$

It is easy to be checked that D is a natural connection with torsion tensor

$$\begin{aligned} T(X, Y, Z) = \frac{1}{4} \{ & \Phi(Y, Z, X) - \Phi(Z, X, Y) - \Phi(Y, \mathcal{P}Z, \mathcal{P}X) + \Phi(\mathcal{P}Z, X, \mathcal{P}Y) \} \\ & + Q'(X, Y, Z) + Q''(X, Y, Z). \end{aligned}$$

From Theorem 2 it follows that $p_2 = Q'$ and $p_4 = Q''$.

Uniqueness. Let D' be another natural connection with torsion tensor T' satisfying the conditions of the theorem. If p'_i ($i=1, 2, 3, 4$) are the components of T' ,

Theorem 3 implies $p'_1 = p_1$, $p'_3 = p_3$. Since $p'_2 = Q' = p_2$, $p'_4 = Q'' = p_4$, then $T' = T$. Hence $D' = D$.

Theorem 4 gives a one-to-one mapping of the set of the natural connections onto the set of pairs (Q', Q'') having the properties (i) and (ii).

Definition 4. A natural connection D with torsion tensor T is said to be canonical if

$$(16) \quad T(X, Y, Z) + T(Y, Z, X) + T(\mathcal{P}X, Y, \mathcal{P}Z) + T(Y, \mathcal{P}Z, \mathcal{P}X) = 0.$$

Taking into account Lemma 5, we find that (16) is equivalent to $T \in \mathcal{T}_1 \oplus \mathcal{T}_3$, i. e. $p_2 = p_4 = 0$. Applying Theorem 4, we check

Theorem 5. There exists a unique canonical connection on a Riemannian almost-product manifold.

From Theorem 4 and (7), (11), (12), (13), (14) it follows that the canonical connection D on a Riemannian almost-product manifold is given by one of the following equalities:

$$g(D_X Y, Z) = g(\nabla_X Y, Z) + \frac{1}{4}\{\Phi(X, Y, Z) - 2\Phi(Z, X, Y) - \Phi(X, \mathcal{P}Y, \mathcal{P}Z)\},$$

$$(17) \quad g(D_X Y, Z) = g(\tilde{\nabla}_X Y, Z) - \frac{1}{4}\{3\Phi(X, Y, Z) + 2\Phi(Z, X, Y) + \Phi(X, \mathcal{P}Y, \mathcal{P}Z)\},$$

$$g(D_X Y, Z) = \frac{1}{2}g(\tilde{\nabla}_X Y + \nabla_X Y, Z) - \frac{1}{4}\{\Phi(Z, X, Y) - \Phi(\mathcal{P}Z, X, \mathcal{P}Y)\}.$$

The torsion tensor T of D is given by

$$(18) \quad T(X, Y, Z) = \frac{1}{4}\{\Phi(Y, Z, X) - \Phi(Z, X, Y) - \Phi(Y, \mathcal{P}Z, \mathcal{P}X) + \Phi(\mathcal{P}Z, X, \mathcal{P}Y)\}.$$

4. Canonical conformal group. In this section we consider the group of transformations of the canonical connections generated by the general conformal transformations of the metric.

Let (M, \mathcal{P}, g) be a Riemannian almost-product manifold. The general conformal transformations of the metric g are defined by

$$(19) \quad \bar{g} = u(\operatorname{ch} v \cdot g + \operatorname{sh} v \cdot \tilde{g}), \quad \tilde{\bar{g}} = u(\operatorname{sh} v \cdot g + \operatorname{ch} v \cdot \tilde{g}),$$

where u, v are differentiable functions on M . By $v=0$ (19) is the usual conformal change of the metric g . The manifold $(M, \mathcal{P}, \bar{g})$ is also a Riemannian almost-product manifold.

We take in consideration only local conformal transformations.

Let (M, \mathcal{P}, g) and $(M, \mathcal{P}, \bar{g})$ be conformally related as in (19). The Levi-Civita connections of \bar{g} and $\tilde{\bar{g}}$ are denoted by $\bar{\nabla}$ and $\tilde{\bar{\nabla}}$ respectively. Applying the formula $2g(\nabla_X Y, Z) = Xg(Y, Z) + Yg(X, Z) - Zg(X, Y) + g([X, Y], Z) - g([Y, Z], X) + g([Z, X], Y)$ to $\bar{\nabla}$ and $\tilde{\bar{\nabla}}$ we find

$$(20) \quad \begin{aligned} 2\bar{g}(\bar{\nabla}_X Y, Z) &= 2u \operatorname{ch} v \cdot g(\nabla_X Y, Z) + 2u \operatorname{sh} v \cdot \tilde{g}(\tilde{\nabla}_X Y, Z) \\ &\quad + X(u \operatorname{ch} v g(Y, Z) + Y(u \operatorname{ch} v) g(X, Z) - Z(u \operatorname{ch} v) g(X, Y) \\ &\quad + X(u \operatorname{sh} v) \tilde{g}(Y, Z) + Y(u \operatorname{sh} v) \tilde{g}(X, Z) - Z(u \operatorname{sh} v) \tilde{g}(X, Y), \end{aligned}$$

and

$$(21) \quad \begin{aligned} \tilde{g}(\tilde{\nabla}_X Y, Z) &= 2u \operatorname{sh} v \cdot g(\nabla_X Y, Z) + 2u \operatorname{ch} v \cdot \tilde{g}(\tilde{\nabla}_X Y, Z) \\ &\quad + X(u \operatorname{sh} v)g(Y, Z) + Y(u \operatorname{sh} v)g(X, Z) - Z(u \operatorname{sh} v)g(X, Y) \\ &\quad + X(u \operatorname{ch} v)\tilde{g}(Y, Z) + Y(u \operatorname{ch} v)\tilde{g}(X, Z) - Z(u \operatorname{ch} v)\tilde{g}(X, Y). \end{aligned}$$

From (6) and (20) we have

$$(22) \quad \begin{aligned} \bar{\nabla}_X Y &= \nabla_X Y - \operatorname{sh}^2 v \cdot \Phi(X, Y) + \frac{1}{2} \operatorname{sh} 2v \cdot \mathcal{P}\Phi(X, Y) + \frac{1}{2} \left\{ \frac{du(X)}{u} Y + \frac{du(Y)}{u} X \right. \\ &\quad \left. + d\mathcal{v}(X)\mathcal{P}Y + d\mathcal{v}(Y)\mathcal{P}X \right\} + \frac{1}{2} g(X, Y) \left\{ -\operatorname{ch}^2 v \frac{\operatorname{grad} u}{u} + \frac{1}{2} \operatorname{sh} 2v \frac{\mathcal{P} \operatorname{grad} u}{u} \right. \\ &\quad \left. - \frac{1}{2} \operatorname{sh} 2v \cdot \operatorname{grad} v + \operatorname{sh}^2 v \cdot \mathcal{P} \operatorname{grad} v \right\} + \frac{1}{2} \tilde{g}(X, Y) \left\{ -\frac{1}{2} \operatorname{sh} 2v \frac{\operatorname{grad} u}{u} \right. \\ &\quad \left. + \operatorname{sh}^2 v \frac{\mathcal{P} \operatorname{grad} u}{u} - \operatorname{ch}^2 v \cdot \operatorname{grad} v + \frac{1}{2} \operatorname{sh} 2v \cdot \mathcal{P} \operatorname{grad} v \right\}. \end{aligned}$$

Analogously, from (6) and (21) we check

$$(23) \quad \begin{aligned} \tilde{\nabla}_X Y &= \tilde{\nabla}_X Y + \operatorname{sh}^2 v \cdot \Phi(X, Y) - \frac{1}{2} \operatorname{sh} 2v \cdot \mathcal{P}\Phi(X, Y) + \frac{1}{2} \left\{ \frac{du(X)}{u} Y + \frac{du(Y)}{u} X \right. \\ &\quad \left. + d\mathcal{v}(X)\mathcal{P}Y + d\mathcal{v}(Y)\mathcal{P}X \right\} + \frac{1}{2} g(X, Y) \left\{ \operatorname{sh}^2 v \frac{\operatorname{grad} u}{u} - \frac{1}{2} \operatorname{sh} 2v \frac{\mathcal{P} \operatorname{grad} u}{u} \right. \\ &\quad \left. + \frac{1}{2} \operatorname{sh} 2v \cdot \operatorname{grad} v - \operatorname{ch}^2 v \cdot \mathcal{P} \operatorname{grad} v \right\} + \frac{1}{2} \tilde{g}(X, Y) \left\{ \frac{1}{2} \operatorname{sh} 2v \frac{\operatorname{grad} u}{u} \right. \\ &\quad \left. - \operatorname{ch}^2 v \frac{\mathcal{P} \operatorname{grad} u}{u} + \operatorname{sh}^2 v \cdot \operatorname{grad} v - \frac{1}{2} \operatorname{sh} 2v \cdot \mathcal{P} \operatorname{grad} v \right\}. \end{aligned}$$

Subtracting (22) from (23) and taking into account the defining condition (6) of Φ , we obtain

$$(23) \quad \begin{aligned} \bar{\Phi}(X, Y) &= \operatorname{ch} 2v \cdot \Phi(X, Y) - \operatorname{sh} 2v \cdot \mathcal{P}\Phi(X, Y) + \frac{1}{2} g(X, Y) \left\{ \operatorname{ch} 2v \frac{\operatorname{grad} u}{u} \right. \\ &\quad \left. - \operatorname{sh} 2v \frac{\mathcal{P} \operatorname{grad} u}{u} + \operatorname{sh} 2v \cdot \operatorname{grad} v - \operatorname{ch} 2v \cdot \mathcal{P} \operatorname{grad} v \right\} + \frac{1}{2} \tilde{g}(X, Y) \left\{ \operatorname{sh} 2v \frac{\operatorname{grad} u}{u} \right. \\ &\quad \left. - \operatorname{ch} 2v \frac{\mathcal{P} \operatorname{grad} u}{u} + \operatorname{ch} 2v \cdot \operatorname{grad} v - \operatorname{sh} 2v \cdot \mathcal{P} \operatorname{grad} v \right\}. \end{aligned}$$

Adding (22) and (23), we get

$$(25) \quad \begin{aligned} \frac{1}{2} (\tilde{\nabla}_X Y + \bar{\nabla}_X Y) &= \frac{1}{2} (\tilde{\nabla}_X Y + \nabla_X Y) + \frac{1}{2} \left\{ \frac{du(X)}{u} Y + \frac{du(Y)}{u} X + d\mathcal{v}(X)\mathcal{P}Y + d\mathcal{v}(Y)\mathcal{P}X \right\} \\ &\quad - \frac{1}{4} g(X, Y) \left\{ \frac{\operatorname{grad} u}{u} + \mathcal{P} \operatorname{grad} v \right\} - \frac{1}{4} \tilde{g}(X, Y) \left\{ \frac{\mathcal{P} \operatorname{grad} u}{u} + \operatorname{grad} v \right\}. \end{aligned}$$

The third equality of (17) and (25) imply

Theorem 6. *The group of the general conformal transformations (19) generates the group of conformal transformations of the canonical connection*

$$(26) \quad \bar{D}_X Y = D_X Y + \frac{1}{2} \left\{ \frac{du(X)}{u} Y + dv(X) \mathcal{P} Y \right\} + \frac{1}{4} \left\{ \frac{du(Y)}{u} + dv(\mathcal{P} Y) \right\} X + \frac{1}{4} \left\{ \frac{du(\mathcal{P} Y)}{u} + dv(Y) \right\} \mathcal{P} X - \frac{1}{4} g(X, Y) \left\{ \frac{\text{grad } u}{u} + \mathcal{P} \text{grad } v \right\} - \frac{1}{4} \tilde{g}(X, Y) \left\{ \frac{\mathcal{P} \text{grad } u}{u} + \text{grad } v \right\}.$$

As a consequence of (26) the torsion tensors \bar{T} and T of \bar{D} and D , respectively, are related as follows:

$$(27) \quad \bar{T}(X, Y) = T(X, Y) - \frac{1}{4} \left\{ \frac{du(Y)}{u} - dv(\mathcal{P} Y) \right\} X + \frac{1}{4} \left\{ \frac{du(X)}{u} - dv(\mathcal{P} X) \right\} Y + \frac{1}{4} \left\{ \frac{du(\mathcal{P} Y)}{u} - dv(Y) \right\} \mathcal{P} X - \frac{1}{4} \left\{ \frac{du(\mathcal{P} X)}{u} - dv(X) \right\} \mathcal{P} Y.$$

The group of the usual conformal transformations of the metric g is characterized by the conditions $u = e^{2f}$, $v = 0$ in (19). In this case the formulas (24), (26) and (27) are reduced to

$$\begin{aligned} \bar{\Phi}(X, Y) &= \Phi(X, Y) + g(X, Y) \text{grad } f - \tilde{g}(X, Y) \mathcal{P} \text{grad } f, \\ \bar{D}_X Y &= D_X Y + df(X) Y + \frac{1}{2} df(Y) X - \frac{1}{2} g(X, Y) \text{grad } f \\ &\quad + \frac{1}{2} df(\mathcal{P} Y) \mathcal{P} X - \frac{1}{2} g(\mathcal{P} X, Y) \mathcal{P} \text{grad } f, \\ \bar{T}(X, Y) &= T(X, Y) - \frac{1}{2} \{ df(Y) X - df(X) Y - df(\mathcal{P} Y) \mathcal{P} X + df(\mathcal{P} X) \mathcal{P} Y \}. \end{aligned}$$

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Received 06. 06. 1989.