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Resonance Mean-Periodic Solutions of Euler Differential Equations

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Abstract. The Euler operator $\delta = t \frac{d}{dt}$ is considered in the space $C = C(\mathbb{R}_+)$ of the continuous functions on $\mathbb{R}_+ = (0, \infty)$. Nonlocal operational calculi for it are developed and used for solving nonlocal Cauchy boundary value problems for Euler differential equations of the form $P(\delta)y = f$ with a polynomial P . A function $f \in C(\mathbb{R}_+)$ is said to be mean-periodic for the Euler operator with respect to the linear functional Φ (or simply Φ -mean-periodic) if $\Phi_\tau\{f(t\tau)\} = 0$ identically on \mathbb{R}_+ . The solution of Euler differential equations in mean-periodic functions for δ with respect to an arbitrary linear functional Φ reduces to non-local homogeneous Cauchy problems.

Denoting the algebraic equivalent of the Euler differential operator δ by S , the solution of an Euler differential equation $P(\delta)y = f$ in Φ -mean-periodic functions reduces to the interpretation of $y = \frac{1}{P(S)}f$ as a function. This is done both in the non-resonance and in the resonance cases.

Keywords: Euler operator, mean-periodic function, resolvent operator, commutant, Euler convolution.

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RESULTS FOR THE NON-RESONANCE CASE

The non-resonance case of solving Euler differential equations in mean-periodic functions and some preliminary definitions and theorems were considered in previous papers and will be briefly presented in this section.

Nonlocal operational calculi for the Euler operator

Let Φ be a non-zero linear functional on $C = C(\mathbb{R}_+)$. The solution of the elementary boundary value problem

$$\delta y - \lambda y = t \frac{dy}{dt} - \lambda y = f(t), \quad \Phi\{y\} = 0,$$

has the form

$$y = L_\lambda f(x) = \frac{t^\lambda}{E(\lambda)} \Phi_\tau \left\{ \tau^\lambda \int_\tau^t \frac{f(\sigma) d\sigma}{\sigma^{\lambda+1}} \right\},$$

where $E(\lambda) = \Phi_\tau\{\tau^\lambda\}$ is the Euler indicatrix of the functional Φ . L_λ is said to be the resolvent operator of δ with respect to Φ .

We begin with the following non-classical convolution related to the Euler operator, which is considered in details in [3], [4], and [6]:

Theorem 1 *Let Φ be a continuous non-zero linear functional on $C(\mathbb{R}_+)$. Then the operation*

$$(f * g)(t) = \Phi_\tau \left\{ \int_\tau^t f\left(\frac{t\tau}{\sigma}\right) g(\sigma) \frac{d\sigma}{\sigma} \right\} \quad (1)$$

is a separately continuous, bilinear, commutative, and associative operation in $C(\mathbb{R}_+)$, such that

$$L_\lambda f(t) = \left\{ \frac{t^\lambda}{E(\lambda)} \right\} * f.$$

Lemma 1 *If $f \in C^1(\mathbb{R}_+)$, then*

$$\delta(f * g) = (\delta f) * g - \Phi\{f\}g.$$

Here we pay attention also to a very useful property of the convolution (1):

Lemma 2 *The convolution given by (1) is such that*

$$\Phi\{f * g\} = 0 \quad (2)$$

for arbitrary $f, g \in C(\mathbb{R}_+)$.

Further, if $\Phi\{1\} \neq 0$, i.e. if $E(0) \neq 0$, then without loss of generality we may assume that the functional Φ satisfies $\Phi\{1\} = 1$ and then the following representation for $L = L_0$ holds:

$$Lf = \{1\} * f.$$

Let the space $C = C(\mathbb{R}_+)$ be considered as a commutative and associative algebra with the convolution $*$ as multiplication. Next, the commutative ring \mathcal{M} of convolution fractions $\frac{f}{g}$ with g being nonzero non-divisor of zero) is defined. Two convolution fractions $\frac{f}{g}$ and $\frac{f_1}{g_1}$ are considered as equal iff $f * g_1 = g * f_1$. The elements of $C(\mathbb{R}_+)$ may be considered as convolutional fractions due to the embedding

$$f \hookrightarrow \frac{f * \{1\}}{\{1\}}.$$

It embeds the ring $(C(\mathbb{R}_+), *)$ into the ring \mathcal{M} of the convolution fractions.

The reciprocal element

$$S = L^{-1}$$

of L in the ring \mathcal{M} is called the *algebraic Euler operator*. Its relation to the ordinary Euler operator δ is given by

Lemma 3 If $f \in C^1(\mathbb{R}_+)$, then

$$\delta f = Sf - \Phi\{f\},$$

where $\Phi\{f\}$ is the “numerical operator” $\frac{\{\Phi\{f\}\}}{\{1\}}$.

By induction it follows that

$$\delta^k f = S^k f - \Phi\{f\}S^{k-1} - \Phi\{\delta f\}S^{k-2} - \dots - \Phi\{\delta^{k-1} f\}. \quad (3)$$

If λ is such that $E(\lambda) = \Phi_\tau(\tau^\lambda) \neq 0$, then (see [8])

$$\frac{1}{S-\lambda} = \left\{ \frac{t^\lambda}{E(\lambda)} \right\} \text{ and } \frac{1}{(S-\lambda)^k} = \left\{ \frac{1}{(k-1)!} \frac{\partial^{k-1}}{\partial \lambda^{k-1}} \left(\frac{t^\lambda}{E(\lambda)} \right) \right\}, \quad k > 1.$$

Nonlocal Cauchy boundary value problems for Euler equations

The general nonlocal Cauchy boundary value problem for the Euler operator $\delta = t \frac{d}{dt}$ has the form

$$P(\delta)y(t) = F(t), \quad \Phi\{\delta^k y\} = \alpha_k, \quad k = 0, 1, 2, \dots, \deg P - 1, \quad (4)$$

where P is a polynomial, Φ is an arbitrary non-zero linear functional, and α_k are real or complex numbers.

Lemma 4 Let none of the zeros of the polynomial $P(\lambda)$ be a zero of the indicatrix $E(\lambda)$, i.e. $\{\lambda : P(\lambda) = 0\} \cap \{\lambda : E(\lambda) = 0\} = \emptyset$. Then $P(S)$ is a non-divisor of zero in \mathcal{M} .

The case, when $P(S)$ is a non-divisor of zero in \mathcal{M} , is called the *non-resonance case*.

In this non-resonance case the operational approach gives the solution simply by substituting (3) in (4) in order to obtain an usual algebraic equation

$$P(S)y = F + Q(S), \quad (5)$$

where $Q(S)$ is a polynomial of S with $\deg Q < \deg P$. It has the formal solution

$$y = \frac{1}{P(S)}F + \frac{Q(S)}{P(S)}.$$

Using the zeros of the polynomial P , the formal quotients $\frac{1}{P(S)}$ and $\frac{Q(S)}{P(S)}$ can be written as sums of elementary fractions. Representing each such fraction as a function and, then using

$$\frac{1}{S-\lambda} = \left\{ \frac{t^\lambda}{E(\lambda)} \right\} \text{ and } \frac{1}{(S-\lambda)^k} = \left\{ \frac{1}{(k-1)!} \frac{\partial^{k-1}}{\partial \lambda^{k-1}} \left(\frac{t^\lambda}{E(\lambda)} \right) \right\}, \quad \text{for } k > 1, \quad (6)$$

one obtains the solution of the nonlocal Cauchy boundary problem in the nonresonance case.

Mean-periodic solutions of Euler differential equations in the non-resonance case

Definition 1 (Schwartz [10], §22) *A function $f \in C(\mathbb{R}_+)$ is said to be mean-periodic for the Euler operator with respect to a linear functional Φ (shortly Φ -mean-periodic, Euler mean-periodic, or simply mean-periodic) if*

$$\Phi_\tau\{f(t\tau)\} = 0$$

identically on \mathbb{R}_+ .

Let us denote by MP_Φ the subset of the mean-periodic functions in $C(\mathbb{R}_+)$ with respect to the functional Φ .

If we are looking for mean-periodic solutions of the Euler equation $P(\delta)y(t) = f(t)$, it is equivalent to a nonlocal Cauchy problem with *homogeneous* boundary value conditions, i.e.

$$P(\delta)y(t) = f(t), \quad \Phi\{\delta^k y\} = 0, \quad k = 0, 1, 2, \dots, \deg P - 1. \quad (7)$$

The following two theorems concerning properties of the mean-periodic functions in the convolutional algebra $(C(\mathbb{R}_+), *)$ are proved by the authors in [3] and [4]:

Theorem 2 *The mean-periodic functions for the Euler operator δ with respect to any non-zero functional $\Phi : C(\mathbb{R}_+) \rightarrow \mathbb{C}$ form an ideal in the convolutional algebra $(C(\mathbb{R}_+), *)$.*

Theorem 3 *If $f \in MP_\Phi$ and $\{\lambda : P(\lambda) = 0\} \cap \{\lambda : E(\lambda) = 0\} = \emptyset$, then $P(\delta)y = f$ has a unique solution $y \in MP_\Phi$ and it has the explicit form*

$$y = G * f,$$

with $G = \{G(t)\} = \frac{1}{P(S)}$ using the representations (6).

OPERATIONAL METHOD FOR MEAN-PERIODIC RESONANCE SOLUTIONS OF EULER DIFFERENTIAL EQUATIONS

Here we give a result for multiple resonance zeros which generalizes the case of simple resonance zeros considered by the authors in [5].

As it was mentioned above, the Φ -mean-periodic resonance solutions of an Euler differential equation have to satisfy the homogeneous Cauchy boundary value conditions:

$$P(\delta)y(t) = f(t), \quad \Phi\{\delta^k y\} = 0, \quad k = 0, 1, 2, \dots, \deg P - 1. \quad (8)$$

If λ_0 is one of the resonance zeros, let us denote by k_0 and \varkappa_0 its multiplicities in the polynomial P and the Euler indicatrix E , respectively, i.e.

$$P(\lambda) = (\lambda - \lambda_0)^{k_0} Q(\lambda), \quad Q(\lambda_0) \neq 0, \quad (9)$$

and

$$E(\lambda_0) = E'(\lambda_0) = \dots = E^{(\varkappa-1)}(\lambda_0) = 0, E^{(\varkappa_0)}(\lambda_0) \neq 0. \quad (10)$$

Now an idea of S. Grozdev [9] will be used for reducing the resonance case to the non-resonance one.

Denote by \tilde{C}_{λ_0} the subalgebra of $(C(\mathbb{R}_+), *)$ consisting of the functions $f \in C(\mathbb{R}_+)$ satisfying the conditions

$$f * t^{\lambda_0} = 0, f * t^{\lambda_0} \ln t = 0, \dots, f * t^{\lambda_0} \ln^{k_0-1} t = 0. \quad (11)$$

Obviously these conditions are necessary for existing of a mean-periodic solution. Our next task is to show that they are also sufficient.

Lemma 5 *Let λ_0 be a resonance zero of the polynomial P . Then, if \tilde{S} denotes the restriction of S to \tilde{C}_{λ_0} , it holds*

$$\frac{1}{(\tilde{S} - \lambda_0)^n} = \left\{ \frac{t^{\lambda_0} A_n(\ln t)}{E^{(\varkappa)}(\lambda_0)} \right\}^*, \quad (12)$$

where $A_n(x)$, $n = 1, 2, \dots$, are an Appell set of polynomials defined recurrently by

$$A_1(\ln t) = \ln^{\varkappa} t, \quad \Phi_t \left\{ t^{\lambda_0} A_n(\ln t) \right\} = 0, \quad A'_{n+1}(x) = A_n(x), \quad x = \ln t. \quad (13)$$

We have to consider two cases: 1) $k_0 \leq \varkappa_0$ and 2) $k_0 > \varkappa_0$.

Case 1: $k_0 \leq \varkappa_0$

For any λ which is not a zero of the polynomial P we can use the representation

$$R_\lambda \tilde{f} = \left\{ \frac{t^\lambda}{E(\lambda)} \right\} * \tilde{f} \quad \text{for } \tilde{f} \in \tilde{C}_{\lambda_0} \text{ and } P(\lambda) \neq 0. \quad (14)$$

Then it is natural to let λ tend to λ_0 :

$$\lim_{\lambda \rightarrow \lambda_0} R_\lambda \tilde{f} = \lim_{\lambda \rightarrow \lambda_0} \left\{ \frac{t^{\lambda_0} t^{\lambda - \lambda_0}}{E(\lambda)} \right\} * \tilde{f}. \quad (15)$$

Representing

$$\begin{aligned} t^{\lambda - \lambda_0} &= e^{(\lambda - \lambda_0) \ln t} = \sum_{k=0}^{\varkappa_0-1} \frac{(\lambda - \lambda_0)^k \ln^k t}{k!} + \\ &+ (\lambda - \lambda_0)^{\varkappa_0} \left[\frac{\ln^{\varkappa_0} t}{\varkappa_0!} + (\lambda - \lambda_0) \sum_{k=\varkappa_0+1}^{\infty} \frac{(\lambda - \lambda_0)^{k-\varkappa_0-1} \ln^k t}{k!} \right] \end{aligned} \quad (16)$$

and

$$E(\lambda) = (\lambda - \lambda_0)^{\varkappa_0} \left[\frac{E^{(\varkappa_0)}(\lambda_0)}{\varkappa_0!} + (\lambda - \lambda_0) \sum_{k=\varkappa_0+1}^{\infty} \frac{(\lambda - \lambda_0)^{k-\varkappa_0-1} E^k(\lambda_0)}{k!} \right], \quad (17)$$

it follows that

$$\lim_{\lambda \rightarrow \lambda_0} R_\lambda \tilde{f} = \left\{ \frac{t^{\lambda_0} \ln^{\varkappa_0} t}{E^{(\varkappa_0)}(\lambda_0)} \right\} * \tilde{f}. \quad (18)$$

Now we can express the elementary fractions needed for the solution as

$$\frac{1}{\tilde{S} - \lambda_0} = \left\{ \frac{t^{\lambda_0} \ln^{\varkappa_0} t}{E^{(\varkappa_0)}(\lambda_0)} \right\}^* \quad (19)$$

and by induction

$$\frac{1}{(\tilde{S} - \lambda_0)^n} = \left\{ \frac{t^{\lambda_0} A_n(\ln t)}{E^{(\varkappa_0)}(\lambda_0)} \right\}^*, \quad (20)$$

where A_n are polynomials of Appel type in the sense of Bourbaki [1] (Chapter 6.2) related to the Euler operator with the properties

$$A_1(\ln t) = \ln^{\varkappa_0} t, \quad \Phi_t \left\{ t^{\lambda_0} A_n(\ln t) \right\} = 0, \quad A'_{n+1}(x) = A_n(x), \quad x = \ln t.$$

Later we will give the form of the solution of (8) and the part corresponding to the resonance zero $\lambda = \lambda_0$ depends in this case on k_0 arbitrary constants.

Case 2: $k_0 > \varkappa_0$

In this case the representation of $\frac{1}{(\tilde{S} - \lambda_0)^n}$ is the same. The only difference is that the part of the solution corresponding to the resonance zero $\lambda = \lambda_0$ depends on \varkappa_0 arbitrary constants as it will be indicated below. \square

General solution

In the case when the polynomial P in (8) has more than one multiple resonance zeros we can proceed as follows:

Let the resonance zeros be $\lambda_1, \dots, \lambda_m$ with multiplicities k_1, \dots, k_m .

In fact, after removing the non-resonance zeros of P , we want to find the mean-periodic solutions w of the auxiliary equation

$$(\delta - \lambda_1)^{k_1} \dots (\delta - \lambda_m)^{k_m} w(t) = F(t). \quad (21)$$

Lemma 5 shows how to treat the case of only one zero. In the general case we can take for each zero λ_μ , $\mu = 1, \dots, m$, the subalgebra \tilde{C}_{λ_μ} of $(C(\mathbb{R}_+), *)$ consisting of the functions $f \in C(\mathbb{R}_+)$ satisfying the conditions

$$f * t^{\lambda_\mu} = 0, f * t^{\lambda_\mu} \ln t = 0, \dots, f * t^{\lambda_\mu} \ln^{k_\mu-1} t = 0. \quad (22)$$

Next, denote by \tilde{L} the restriction of L to the intersection $\tilde{C}_{\lambda_1, \dots, \lambda_m} = \cap_{\mu=1}^m \tilde{C}_{\lambda_\mu}$. Let $\tilde{S} = \frac{1}{\tilde{L}}$. In the ring $\tilde{M}_{\lambda_1, \dots, \lambda_m}$ of the convolutional fractions in $\tilde{C}_{\lambda_1, \dots, \lambda_m}$, equation (21) takes the form

$$(\tilde{S} - \lambda_1)^{k_1} \dots (\tilde{S} - \lambda_m)^{k_m} w = F. \quad (23)$$

The elements $(\tilde{S} - \lambda_k)$, $k = 1, 2, \dots, m$, are non-divisors of zero in $\tilde{M}_{\lambda_1, \dots, \lambda_m}$ and the formal solution of (23) is

$$w = \frac{1}{(\tilde{S} - \lambda_1)^{k_1} \dots (\tilde{S} - \lambda_m)^{k_m}} F. \quad (24)$$

Next step is to represent the algebraic multiplier $\frac{1}{(\tilde{S} - \lambda_1)^{k_1} \dots (\tilde{S} - \lambda_m)^{k_m}}$ in the form

$$\sum_{\mu=1}^m \sum_{j=1}^{k_\mu} \frac{B_{\mu,j}}{(\tilde{S} - \lambda_\mu)^j}.$$

The general solution of the auxiliary equation (21), which contains only the resonance zeros of P from the Euler equation (8), is then a sum

$$w = \sum_{\mu=1}^m \left(\sum_{j=1}^{k_\mu} B_{\mu,j} \left\{ \frac{t^{\lambda_\mu} A_{\mu,j}(\ln t)}{E^{(\lambda_\mu)}(\lambda_\mu)} \right\} * F + \sum_{j=1}^{\min\{k_\mu, \lambda_\mu\}} C_{\mu,j} t^{\lambda_\mu} \ln^{j-1} t \right)$$

with a uniquely determined constants $B_{\mu,j}$ and arbitrary constants $C_{\mu,j}$, $\mu = 1, \dots, m$, $j = 1, \dots, \min\{k_\mu, \lambda_\mu\}$. Here $A_{\mu,j}$, $\mu = 1, \dots, m$, $s = 1, 2, \dots$, are m polynomial systems of Appel type in the sense of Bourbaki [1] (Chapter 6.2) related to the Euler operator with the properties

$$A_{\mu,1}(\ln t) = \ln^{\lambda_\mu} t, \quad \Phi_t \left\{ t^{\lambda_\mu} A_{\mu,s}(\ln t) \right\} = 0, \quad A'_{\mu,s+1}(x) = A_{\mu,s}(x), \quad x = \ln t.$$

In order to have a full solution of (8), the non-resonance zeros have also to be taken into account using for them the representations (6).

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