CB processes with quadratic competition and conditioning on the non-extinction

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Let $d \geq 0, \ c > 0,$ and $\Pi = \left(\Pi_j, j \in \mathbb{Z}^+\right)$ such that $0 < \rho := \sum_{i \in \mathbb{Z}^+} \Pi_i < \infty.$

We will consider a continuous time random Markov chain

$$(L_t, t \geq 0)$$

that takes values in $\ensuremath{\mathbb{Z}}^+,$ whose infinitesimal generator is given by

$$Q_{i,j} = \begin{cases} i\Pi_{j-i}, & j > i, \& i \ge 1; \\ di + c\binom{i}{2}, & i \ge 1, j = i - 1; \\ -i\left(d + \rho + \frac{c(i-1)}{2}\right), & j = i. \end{cases}$$

Thus, we start with *i*-individuals equipped with independent exponential clocks with parameter $\rho + d$, and $\binom{i}{2}$ exponential clocks of parameter c. When either of these clocks ring, if it is of the first type, then the corresponding particle begets a random number of descendants according to a probability proportional to Π , or dies with a probability proportional to d; while if the clock is of the second type then a pair of individuals among the $\binom{i}{2}$ in the population compete and only one of them survives.

 This model was introduced and studied by A. Lambert under the assumption of a finite log moment

$$\sum_{i\geq 2}\log(i)\Pi_i<\infty$$

.

- He proved that if d = 0, the process L is positive recurrent.
- While if d > 0, then the process hits zero in a finite time a.s. In that case, the process comes down from infinity.
- Under some assumptions, Lambert proved that such processes admit a renormalization of time and space such that the scaling limit is a weak solution to the stochastic differential equation of the type

$$dZ_t = bZ_t dt + \sqrt{\gamma Z_t} dB_t - cZ_t^2 dt;$$

with $b, \gamma \geq 0$, some parameters.

Logistic continuous state branching process

Let $(Z_t, t \ge 0)$ be the size of a population evolving in continuous time and space along the dynamics :

- **branching**: each individual reproduces or dies independently with a same law (classical *CB*'s dynamics).
- quadratic competition : pairwise fights at constant rate $c \ge 0$ (quadratic negative drift).

$$\mathrm{d}Z_t = \text{ } \ll \text{CB dynamics } \gg -\frac{c}{2}Z_t^2\mathrm{d}t.$$

The competition <u>breaks</u> the branching property.

The process Z has been introduced by Lambert (2005) and is called *logistic* CB process (LCB).

Let Ψ be a branching mechanism :

$$\Psi(x) := \frac{\sigma^2}{2} x^2 - \gamma x + \int_0^\infty \left(e^{-xy} - 1 + xy \mathbb{1}_{\{y \le 1\}} \right) \pi(\mathrm{d}y) \tag{1}$$

with $\sigma \geq 0, \gamma \in \mathbb{R}$ and π a Lévy measure.

Definition/Theorem

A LCB(Ψ , c) is solution to the stochastic equation :

$$Z_{t} = z + \sigma \int_{0}^{t} \sqrt{Z_{s}} dB_{s} + \gamma \int_{0}^{t} Z_{s} ds + \int_{0}^{t} \int_{0}^{Z_{s}} \int_{1}^{\infty} y \mathcal{M}(ds, du, dy)$$
$$+ \int_{0}^{t} \int_{0}^{Z_{s}} \int_{0}^{1} y \bar{\mathcal{M}}(ds, du, dy) - \frac{c}{2} \int_{0}^{t} Z_{s}^{2} ds,$$
(2)

with B a Brownian motion, \mathcal{M} an indep. PRM with intensity $\mathrm{d} \mathrm{s} \mathrm{d} u \pi(\mathrm{d} y)$ and $\bar{\mathcal{M}}(\mathrm{d} s, \mathrm{d} u, \mathrm{d} y) := \mathcal{M}(\mathrm{d} s, \mathrm{d} u, \mathrm{d} y) - \mathrm{d} \mathrm{s} \mathrm{d} u \pi(\mathrm{d} y)$.

Aim: Given a LCB process Z satisfying

$$\mathbb{P}_{\mathbf{z}}(Z_t \underset{t \to \infty}{\longrightarrow} 0) = 1$$
 (almost-sure asymptotic extinction),

we wish to condition the process on the negligible event

$$\mathscr{S}:=\{Z_t\underset{t\to\infty}{\longrightarrow}0\}^c.$$

It depends in general on how the event ${\mathscr S}$ is approached.

(1) When there is extinction in finite time a.s., i.e.

$$\{Z_t \underset{t\to\infty}{\longrightarrow} 0\} = \{\zeta_0 < \infty\},\$$

with $\zeta_0 := \inf\{t > 0 : Z_t = 0\}$, we could seek a conditioning « along the approximation » :

$$\bigcap_{s>0} \{\zeta_0>s\}.$$

This is the notion of Q-process:

$$\mathbb{Q}_z(\Lambda) = \lim_{s \to \infty} \mathbb{P}_z(\Lambda | \zeta_0 > t + s), \forall \Lambda \in \mathcal{F}_t.$$

(2) We could also try to force the process to go above any levels before being close to 0, i.e. we approach $\mathcal S$ by :

$$\bigcap_{b>0} \{\zeta_b^+ < \zeta_0\}, \text{ with } \zeta_b^+ := \inf\{t > 0 : Z_t > b\}.$$

Those methods do not seem to apply to LCBs without strong assumptions on Ψ . We will approach $\mathscr S$ in a different way.



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Extinction and total progeny

We call total progeny,

$$J:=\int_0^\infty Z_t\mathrm{d}t.$$

Proposition (F., Rivero, Winter 24 (c>0), Bingham 75 (c=0))

The following equality holds a.s.

$$\{Z_t \underset{t \to \infty}{\longrightarrow} 0\} = \{J < \infty\}.$$

$$\bullet \ \ J < \infty \ \mathbb{P}_{z} \text{-p.s. iff} \ \begin{cases} \Psi'(0+) \geq 0 & \text{if } c = 0, \\ \mathbb{H} & \text{if } c > 0. \end{cases}$$

$$\bullet \ \mathbb{E}_z(J) < \infty \ \text{iff} \ \begin{cases} \Psi'(0+) > 0 & \text{if } c = 0, \\ \Psi(\infty) = \infty \ \& \ \int^\infty \log y \ \pi(\mathrm{d}y) < \infty & \text{if } c > 0. \end{cases}$$

We will approach the survival event ${\mathscr S}$ by forcing the total progeny to be infinite, in the following way

$$\mathscr{S} = \bigcap_{\theta > 0} \{ J > e/\theta \} \tag{3}$$

with e an exponential r.v. of parameter 1 independent from Z.

Construction of LCB and proof of the proposition

The generator of the LCB $\it Z$ takes the form :

$$\mathscr{L}f(z) := z \mathrm{L}^{\Psi}f(z) - \frac{c}{2}z^2 f'(z)$$

with L^{Ψ} the generator of a Lévy process Y of Laplace exponent Ψ .

Factorization:

$$\mathscr{L}f(z) = z\left(\mathrm{L}^{\Psi}f(z) - \frac{c}{2}zf'(z)\right) =: z\mathscr{G}f(z)$$

Let $J_s:=\int_0^s Z_u \mathrm{d}u$, $J_\infty=J$ and the random clock :

$$C_t := \inf\{s \ge 0 : J_s > t\}$$

Lamperti's Transformation: The time-changed process

$$(R_t := Z_{C_t}, t \leq J_{\infty})$$

is a positive Markov process with generator \mathscr{G} , it satisfies

$$\mathrm{d}R_t = \mathrm{d}Y_t - \frac{c}{2}R_t\mathrm{d}t, t \le \sigma_0$$

where $\sigma_0 := \inf\{t > 0 : R_t = 0\}.$

ightarrow By the time change, $Z_t=R_{J_t}, \forall t\geq 0$, $\sigma_0=J_{\infty}=J$, and

$${J<\infty} = {\sigma_0<\infty} = {Z_t \xrightarrow[t\to\infty]{} 0}.$$

Asymptotic extinction condition :

$$\mathbb{H}: \Psi(\infty) = \infty \text{ and } \mathcal{E}:= \int_0^{x_0} \frac{1}{u} e^{\int_u^{x_0} \frac{2\Psi(v)}{cv} \mathrm{d}v} \mathrm{d}u = \infty.$$

Examples satisfying $\mathbb H$

• Stable and Neveu mechanisms :

$$\Psi(x) := ax^{\alpha} - \gamma x, \ \forall x \ge 0, \ \text{for} \ \alpha \in (1, 2], \gamma \in \mathbb{R}, a > 0,$$

$$\Psi(x) := x \log x, \ \forall x \ge 0$$

• Let $\alpha \in (0, c/2], a > 0, \beta \in [1, 2]$ and Ψ such that

$$\Psi(x) \underset{x \to 0}{\sim} -\alpha/\log(1/x)$$
 and $\Psi(x) \underset{x \to \infty}{\sim} ax^{\beta}$.

NB : here
$$\int_0^1 \frac{|\Psi(x)|}{x} dx = \infty$$
.

Fact:

$$\int_0^\infty \frac{\Psi(x)}{x} dx > -\infty \iff \int_0^\infty \log y \pi(dy) < \infty \implies \mathcal{E} = \infty.$$

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Without competition

Theorem (Foucart, R., Winter (2024))

Let c=0 and Ψ (sub)-critical : $\varrho:=\Psi'(0+)\geq 0$ ($\Longleftrightarrow J<\infty$ \mathbb{P}_z -a.s.).

• The function h(z) = z is **excessive** and $\forall z > 0$,

$$\mathbb{P}_z^{\uparrow}(\Lambda,t<\zeta):=\mathbb{E}_z\left(\frac{Z_t}{z}\mathbb{1}_{\Lambda}\right)=\lim_{\theta\to 0}\mathbb{P}_z\big(\Lambda,J_t\leq \mathrm{e}/\theta\,\big|J\geq \mathrm{e}/\theta\big)\,.$$

- If $\varrho=0$, $\mathbb{E}_z(J)=\infty$, Z martingale, $\zeta=\infty$, \mathbb{P}_z^{\uparrow} -a.s.
- If $\varrho > 0$, $\mathbb{E}_z(J) < \infty$, Z supermartingale, $\zeta < \infty$, \mathbb{P}_z^{\uparrow} -a.s.
- $(Z, \mathbb{P}_z^{\uparrow})$ satisfies

$$\begin{split} Z_t = & z + \ll \ \textit{CB}(\Psi) \ \textit{dynamics} \ \gg \\ & + \sigma^2 t + \int_0^t \int_0^\infty y \mathcal{I}(\mathrm{d} s, \mathrm{d} y), \ t < \zeta, \end{split}$$

with \mathcal{I} a Poisson mes. of intensity ds $y\pi(dy)$ and $\zeta \stackrel{\text{Law}}{=} \operatorname{Exp}(\varrho)$.



The additional term

$$\left(\sigma^2 t + \int_0^t \int_0^\infty y \mathcal{I}(\mathrm{d} s, \mathrm{d} y), t \ge 0\right)$$

is a subordinator of Laplace exponent Ψ' .

- ightarrow immigration dynamics independent on the population size,
- ightarrow The process $(Z, \mathbb{P}_z^{\uparrow})$ is a CBI with mechanisms Ψ and Ψ' .
- If there is extinction in finite time, (NASC : $\int_{-\infty}^{\infty} \frac{\mathrm{d}u}{\Psi(u)} < \infty$),

$$(Z,\mathbb{P}_z^{\uparrow})\stackrel{\mathrm{Law}}{=} Q$$
-process **killed** at an indep. time $\sim \mathrm{Exp}(\varrho)$

• The process $(Z, \mathbb{P}_z^{\uparrow})$ starts from z = 0 and 0 is interpreted as an *immortal* individual.

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With competition: looking for an excessive function

Theorem (Foucart, R., Winter (2024))

Under the hypothesis \mathbb{H} and c > 0, there exists a constant c_{θ} such that

$$c_{\theta}\mathbb{P}_{z}(J\geq e/\theta),$$

has a limit as $\theta \to 0+$, for any z > 0;

$$\lim_{\theta \to 0} c_{\theta} \mathbb{P}_{z}(J \geq e/\theta) =: h(z) = \int_{0}^{\infty} (1 - e^{-xz}) \underbrace{\frac{1}{x} e^{-\int_{x_{0}}^{x} \frac{2\Psi(u)}{cu} du}}_{=:s(dx)} dx.$$

Moreover, when

$$\Psi(\infty) = \infty \& \int_{-\infty}^{\infty} \log y \, \pi(\mathrm{d}y) < \infty,$$

we have $h(z) = \mathbb{E}_z(J) < \infty, z > 0$.



Theorem (Foucart, R., Winter (2024))

Assume \mathbb{H} and c>0. Let $x_0>0$ fixed. Set $\forall z\in [0,\infty),\ h(z):=\int_{c}^{\infty}(1-e^{-xz})\frac{1}{e^{-\int_{x_0}^{x}\frac{2\Psi(u)}{cu}\,\mathrm{d}u}}\mathrm{d}x\,.$

- h is of Bernstein form \Longrightarrow positive \uparrow , $C^{\infty}(0,\infty)$,
 - h(0) = 0, $h(\infty) = \infty$, $h'(0) < \infty$ and $\int_{-\infty}^{\infty} h(y)\pi(\mathrm{d}y) < \infty$.
- $2 \forall z \geq 0,$

$$\mathscr{L}h(z)=-\frac{c\ell}{2}z\leq 0$$

with
$$\ell := \exp\left(\int_0^{x_0} \frac{2\Psi(u)}{cu} du\right) \ge 0$$
, and

$$\ell > 0$$
 iff $\int_{-\infty}^{\infty} \log y \, \pi(\mathrm{d}y) < \infty$.



Study of the *h*-transformed process

 $ig(h(Z_t), t \geq 0ig)$ is a \mathbb{P}_z -supermartingale and we define :

$$\mathbb{1}_{\{t<\zeta\}}\mathrm{d}\mathbb{P}_z^\uparrow:=\frac{h(Z_t)}{h(z)}\mathrm{d}\mathbb{P}_z,\ \ \text{on}\ \mathcal{F}_t,\ \forall t\geq 0\ \text{and}\ z>0,$$

with ζ the lifetime of $(Z, \mathbb{P}_z^{\uparrow})$ and ∞ is the cemetery state.

Theorem (Foucart, R., Winter (2024))

$$\mathbb{P}_{z}^{\uparrow}(\Lambda, t < \zeta) = \lim_{\theta \to 0} \mathbb{P}_{z}(\Lambda, J_{t} \leq e/\theta \, \big| \, J \geq e/\theta \big) \,, \forall \, \Lambda \in \mathcal{F}_{t}, \forall t \geq 0$$

② $(Z, \mathbb{P}_z^{\uparrow})$ is a $(0, \infty]$ -valued Feller process, $\zeta < \infty$ \mathbb{P}_z^{\uparrow} -a.s., and

3

$$\mathbb{P}_{z}^{\uparrow}\left(\inf_{0\leq s<\zeta}Z_{s}\leq a\right)=\frac{h(a)}{h(z)},\ \forall z>a\geq 0.$$

In particular, $\inf_{0 \leq t < \zeta} Z_t > 0$, \mathbb{P}_z^{\uparrow} -a.s. for all z > 0.

For all z > 0 and y > 0, let

$$b(z) := z \frac{h'(z)}{h(z)}, \ q(z,y) := \frac{z}{h(z)} (h(z+y) - h(z)) \ \text{and} \ k(z) := \frac{c\ell}{2} \frac{z}{h(z)}.$$

Theorem (Foucart, R., Winter (2024))

 $(Z, \mathbb{P}_z^{\uparrow})$ has same law as the weak solution of the stochastic equation below, killed at time $\zeta := \inf\{t > 0 : \int_0^t \frac{k}{\zeta}(Z_s) \mathrm{d}s \ge e\}.$

$$\begin{split} Z_t = & z + \ll \ \textit{LCB}(\Psi, c) \ \textit{dynamics} \gg \\ & + \sigma^2 \int_0^t \frac{\textit{b}(Z_s) \mathrm{d}s}{\mathsf{b}} + \int_0^t \int_0^{\textit{q}(Z_{s-}, y)} \int_0^\infty \textit{y} \mathcal{N}(\mathrm{d}s, \mathrm{d}\textit{u}, \mathrm{d}\textit{y}), \ t < \zeta, \end{split}$$

where e is a standard exponential r.v., $\mathcal N$ a Poisson measure of intensity $\operatorname{ds}\operatorname{du}\pi(\operatorname{dy})$, everything is mutually indep.

 \rightarrow size-dependent immigration, see also Z. Li's work (2019).



Starting from zero : immortal individual

Theorem (Foucart, R., Winter (2024))

•

$$\mathbb{P}_{z}^{\uparrow} \underset{z \to 0+}{\Longrightarrow} \mathbb{P}_{0}^{\uparrow}$$
, in Skorokhod's sense,

with
$$\mathbb{P}_0^{\uparrow}$$
 s.t. $\mathbb{P}_0^{\uparrow}\big(Z_0=0, \exists \, t>0: \forall s\geq t, Z_s>0\big)=1.$

• $(Z, \mathbb{P}_0^{\uparrow})$ is weak solution to the SDE with z=0, where b, q, k are defined at 0 by :

$$b(0) := 1, \quad \forall y > 0, \ q(0, y) := \frac{h(y)}{h'(0)}, \ \text{and} \ k(0) := \frac{c\ell}{2} \frac{1}{h'(0)}.$$

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Study of (Z, \mathbb{P}_z) and $(Z, \mathbb{P}_z^{\uparrow})$

We use **two duality relationships**.

$$\begin{split} &(Z,\mathscr{L}) \overset{\mathsf{Laplace\ dual}}{\longleftrightarrow} (U,\mathscr{A}) \overset{\mathsf{Siegmund\ dual}}{\longleftrightarrow} (V,\mathscr{G}) \\ &\mathbb{E}_z(e^{-xZ_t}) = \mathbb{E}_x(e^{-zU_t}) \text{ and } & \mathbb{P}_x(U_t > y) = \mathbb{P}_y(x > V_t) \end{split}$$

with U et V diffusions, weak solutions to

$$dU_t = \sqrt{cU_t}dB_t - \Psi(U_t)dt, \ U_0 = x$$

$$dV_t = \sqrt{cV_t}dB_t + (c/2 + \Psi(V_t))dt, \ V_0 = y.$$

We call V the **bidual** process of Z.

$$\mathbb{H}$$
 "=" NASC (Feller's tests) for $V_t \underset{t \to \infty}{\longrightarrow} \infty$ a.s.

By combining the dualities, we get

$$\mathbb{E}_{z}(e^{-xZ_{t}}) = \int_{0}^{\infty} z e^{-zy} \mathbb{P}_{y}(V_{t} > x) dy.$$

Recall the excessive function h and $s(\mathrm{d}x) = \frac{1}{x}e^{-\int_x^{\infty}\frac{2\Psi(u)}{cu}\mathrm{d}u}$. The scale function of V vanishing at ∞ is $S(y) := \int_{-\infty}^{\infty} s(\mathrm{d}x)$ and

$$\mathit{h}(z) = \int_0^\infty z \mathrm{e}^{-zy} \mathit{S}(y) \mathrm{d}y \text{ and } \mathbb{E}_z \big(\mathit{h}(\mathit{Z}_t) \big) = \int_0^\infty y \mathrm{e}^{-zy} \mathbb{E}_y \big(\mathit{S}(\mathit{V}_t) \big) \mathrm{d}y.$$

Lemma (Foucart, R., Winter (2024))

 $(h(Z_t), t \geq 0)$ under \mathbb{P}_z ,

• is a strict supermartingale (i.e. this is not a local martingale) when

$$\int_{-\infty}^{\infty} \log y \, \pi(\mathrm{d}y) < \infty \, (\iff \ell > 0 \iff \mathbb{E}_{z}(J) < \infty),$$

• is a strict local martingale (i.e. this is not a martingale) when

$$\int^{\infty} \log y \, \pi(\mathrm{d}y) = \infty \ (\Longleftrightarrow \ell = 0 \Longleftrightarrow \mathbb{E}_{z}(J) = \infty).$$

About *h*-transforms and locally harmonic functions

If T is an $(\mathcal{F}_t)_{t\geq 0}$ -stopping time then for all $A\in\mathcal{F}_T$ et $z\in(0,\infty)$,

$$\mathbb{P}_{z}^{\uparrow}(A, T < \zeta) = \frac{1}{h(z)} \mathbb{E}_{z} \left(h(Z_{T}) \mathbb{1}_{A} \right). \tag{4}$$

Three different situations:

- If $(h(Z_t), t \ge 0)$ is a \mathbb{P}_z -martingale, then $(Z, \mathbb{P}_z^{\uparrow})$ has an infinite lifetime : $\zeta = \infty$, \mathbb{P}_z^{\uparrow} -a.s.
- ② If $(h(Z_t), t \ge 0)$ is a \mathbb{P}_{z^-} strict supermartingale (i.e. this is not a local martingale), $(Z, \mathbb{P}_z^{\uparrow})$ has a finite lifetime and it is killed with positive probability. One has $\mathbb{P}_z^{\uparrow}(Z_{\zeta^-} < \infty) > 0$.
- If $(h(Z_t), t \ge 0)$ is a \mathbb{P}_{z^-} strict local martingale (i.e. this is not a true martingale), $(Z, \mathbb{P}_z^{\uparrow})$ has a finite lifetime but is not killed. It explodes : $Z_{\zeta-} = \infty$, \mathbb{P}_z^{\uparrow} -a.s..

Thank you for you attention!